

# Phase-Change Memory Devices: Fundamentals and Applications (Part I)

Abu Sebastian  
Principal Research Staff Member  
IBM Research - Zurich



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# Acknowledgements

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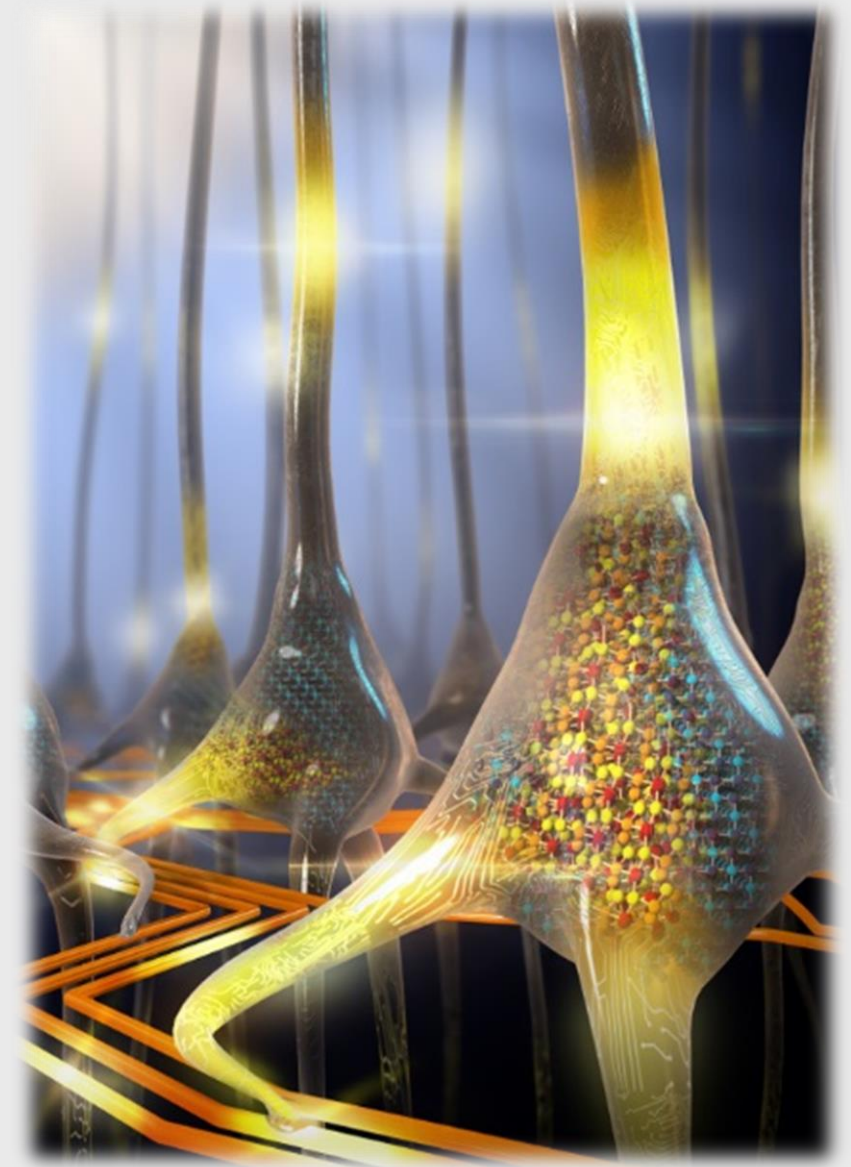


FONDS NATIONAL SUISSE  
SCHWEIZERISCHER NATIONALFONDS  
FONDO NAZIONALE SVIZZERO  
SWISS NATIONAL SCIENCE FOUNDATION



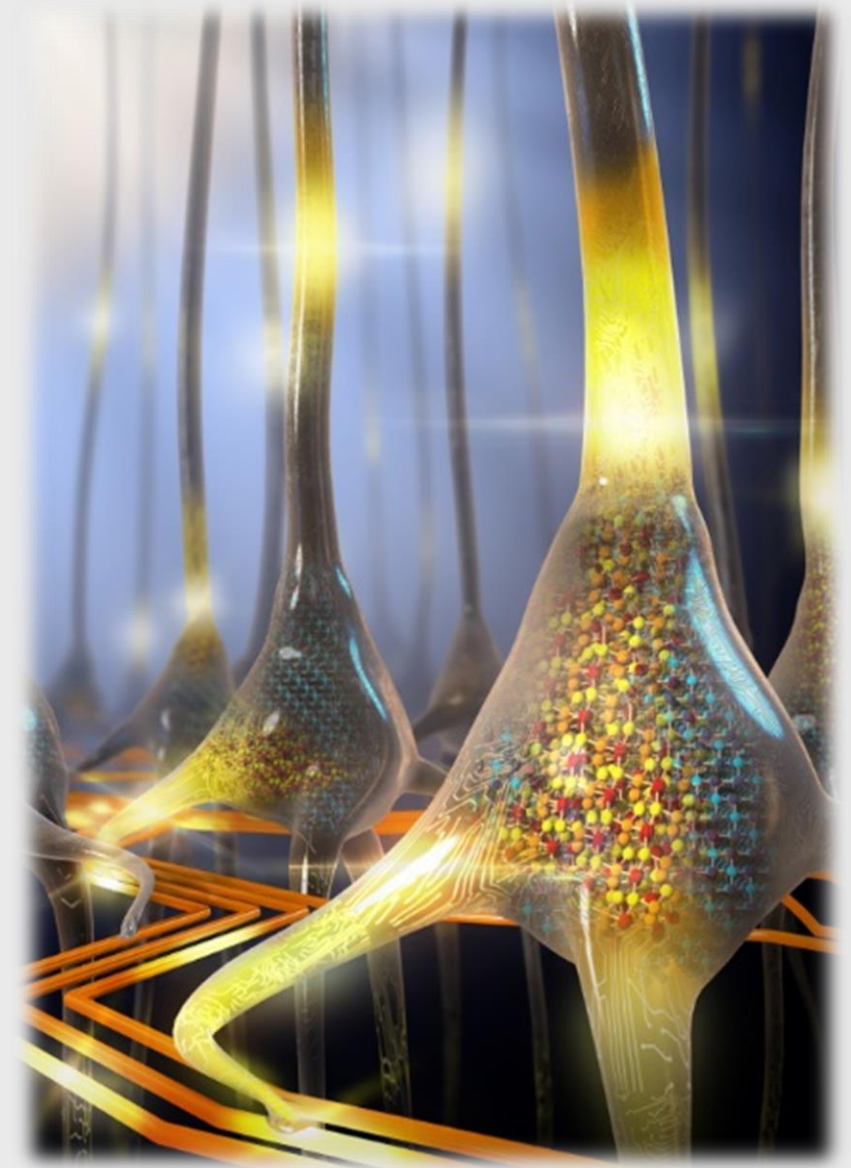
# Outline: Part I

- **Introduction to PCM**
- **PCM device physics**
  - ✓ **Electrical system**
    - Subthreshold electrical transport
    - Threshold switching
  - ✓ **Thermal system**
  - ✓ **Structural dynamics**
    - Melt-quench process
    - Crystallization
    - Structural relaxation
- **Key challenges and device-level advances**
  - ✓ Projected PCM
  - ✓ Single-elemental PCM



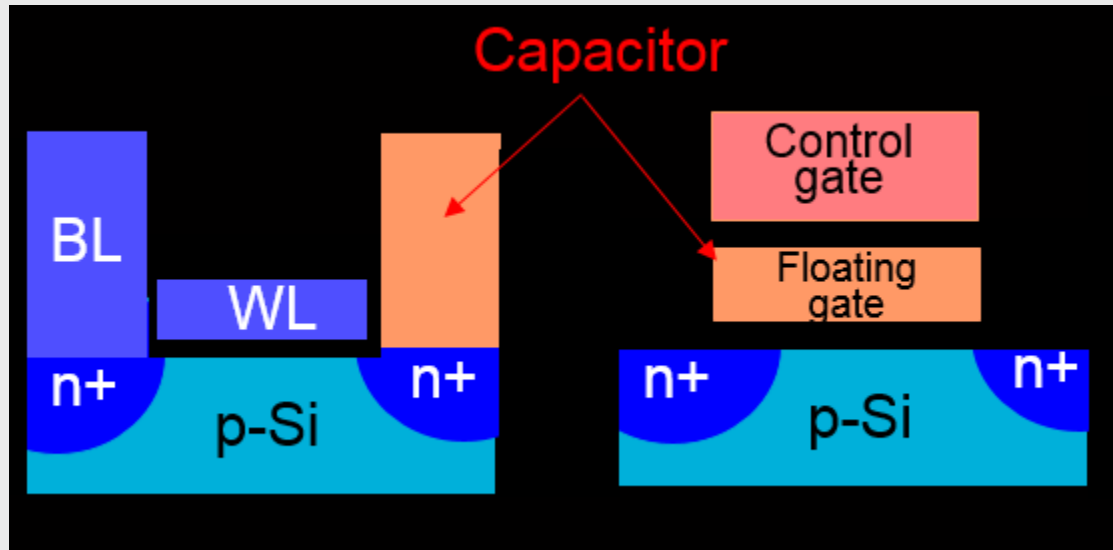
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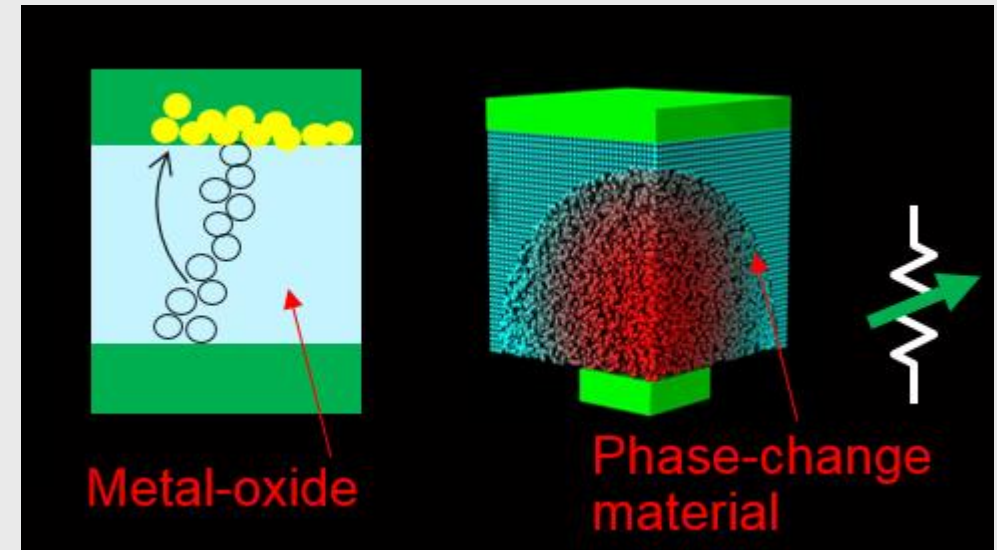


# Resistive memory devices

“Charge on a capacitor”



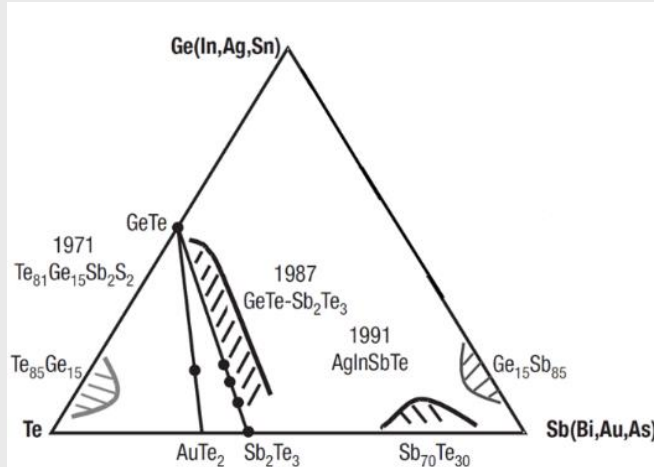
“Alternate atomic arrangements”



- Difference in atomic arrangements induced by the application of electrical pulses and measured as a difference in electrical resistance
- **Resistive memory devices** or “**memristive**” devices
- Based on physical mechanisms such as **ionic drift** and **phase transition**

# Phase-change memory

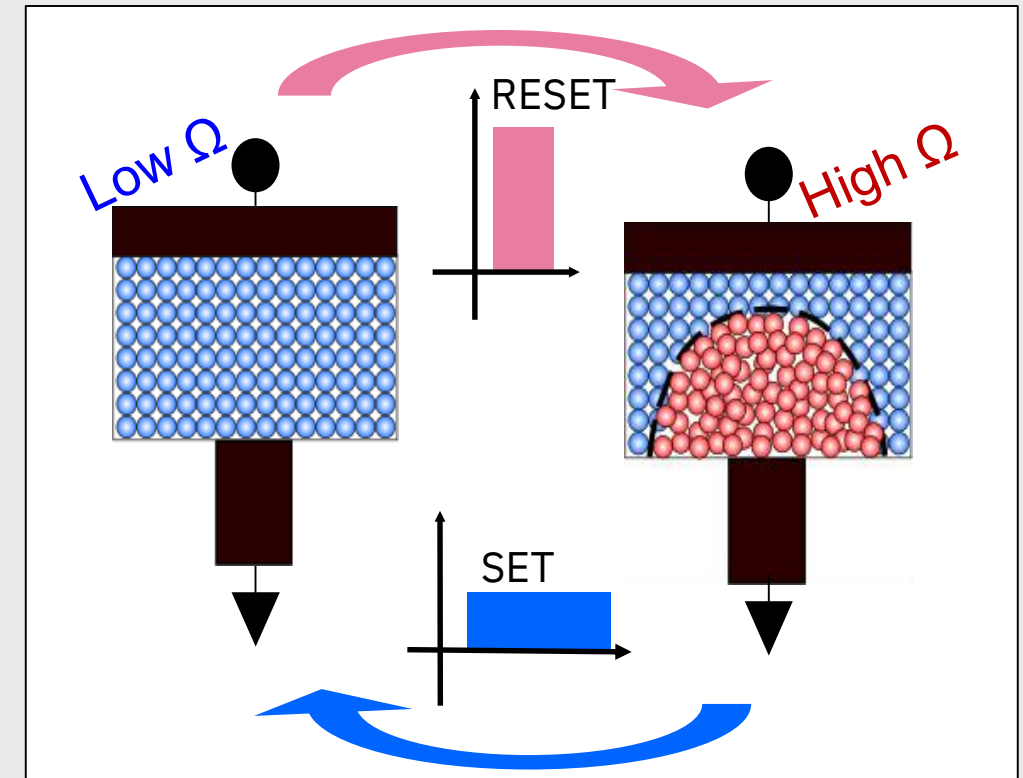
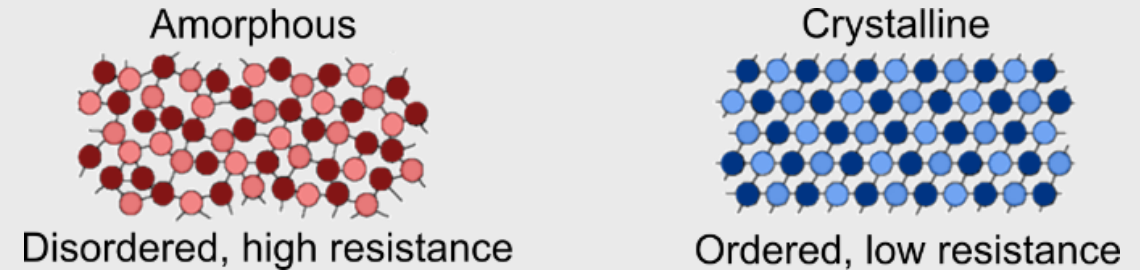
## Commonly used phase change materials



*Wuttig & Yamada, Nature Materials, 2007*

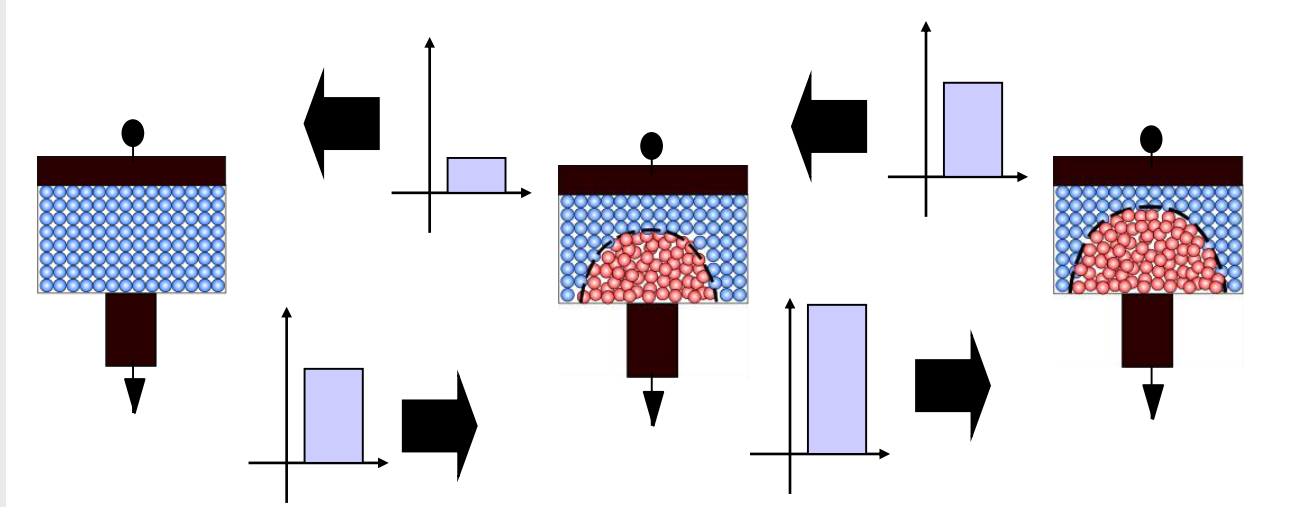
*Burr et al., JETCAS, 2016*

- A nanometric volume of phase change material between two electrodes
- “WRITE” Process
  - ✓ By applying a voltage pulse the material can be changed from crystalline phase (SET) to amorphous phase (RESET)
- “READ” process
  - ✓ Low-field electrical resistance

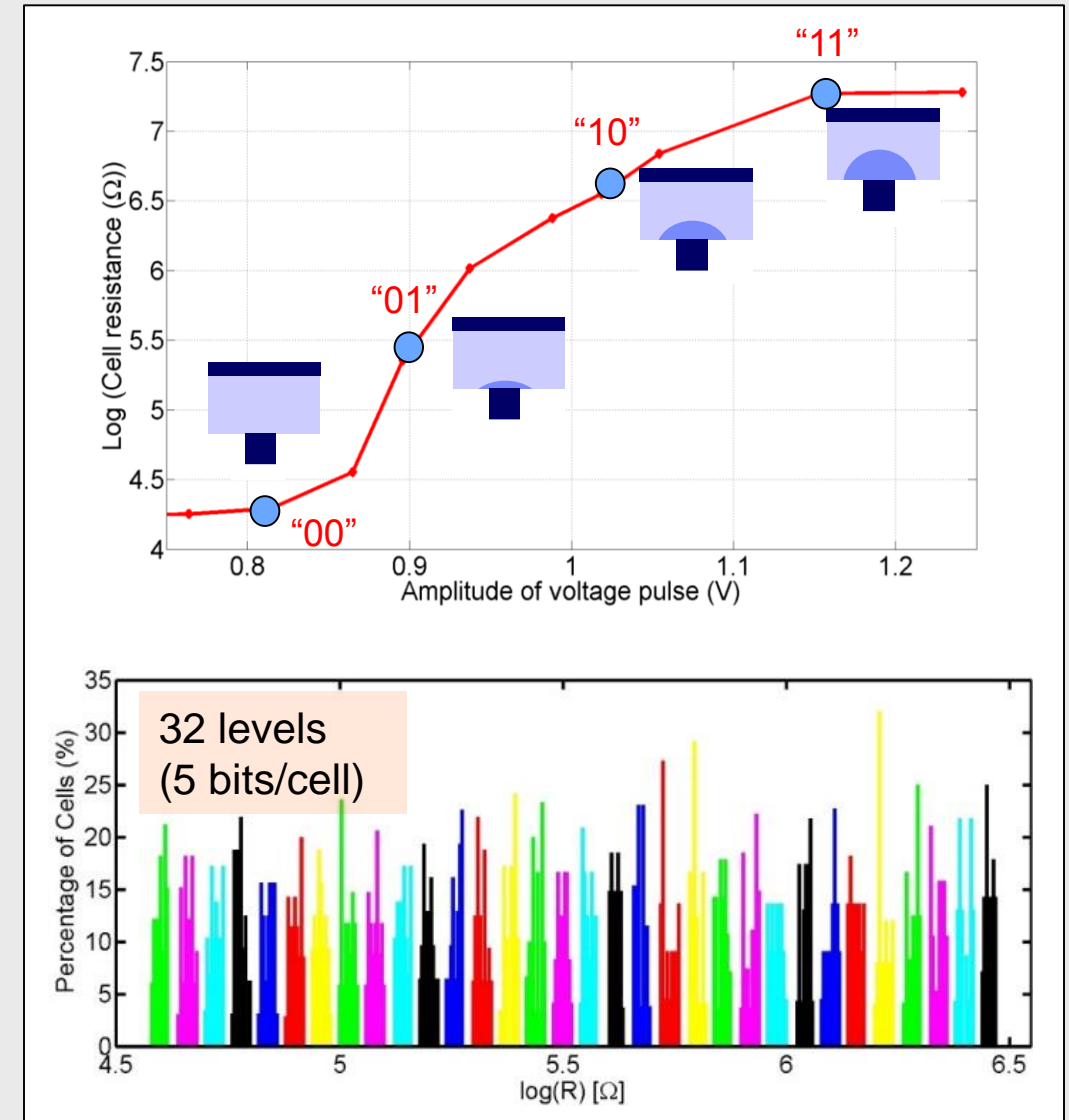




# Multi-level storage



- It is possible to continuously vary the crystalline/amorphous phase configuration and hence the resulting electrical resistance
- Can store more than 1 bit of information in a cell

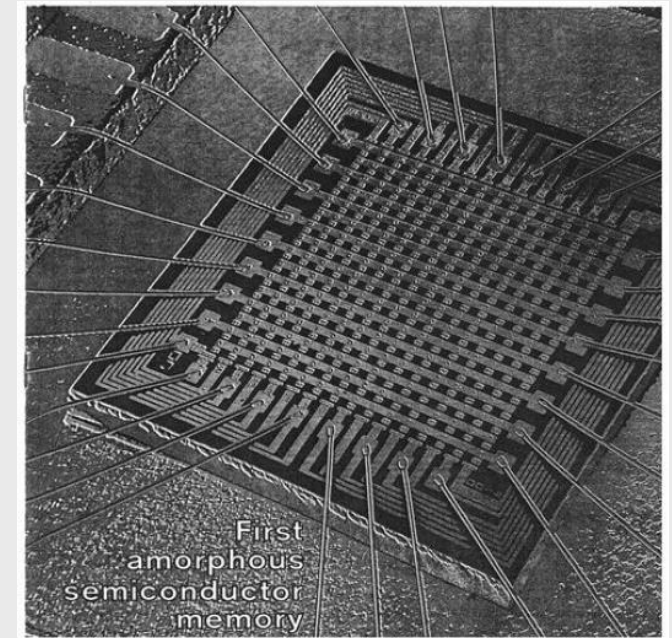
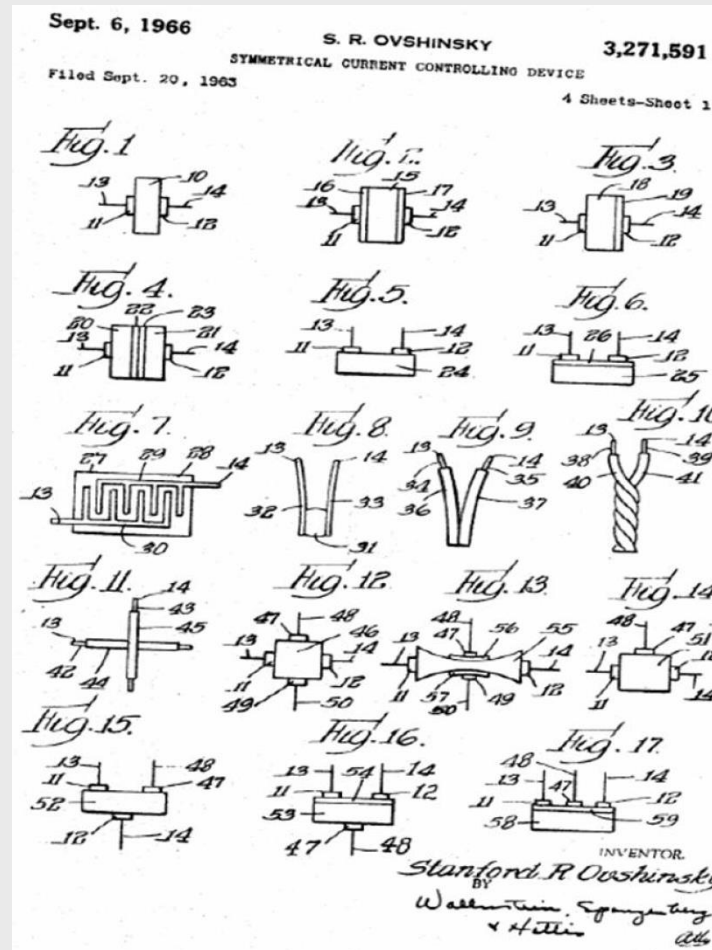


*Papandreou et al., ISCAS, 2011*

# A brief history of phase change memory



**Stan Ovshinsky (1960s)**



**R. G. Neale, D. L. Nelson and G. E. Moore., Electronics, 1970**

Capacity: 256 bits

RESET: ~200mA, <25V, 5 us

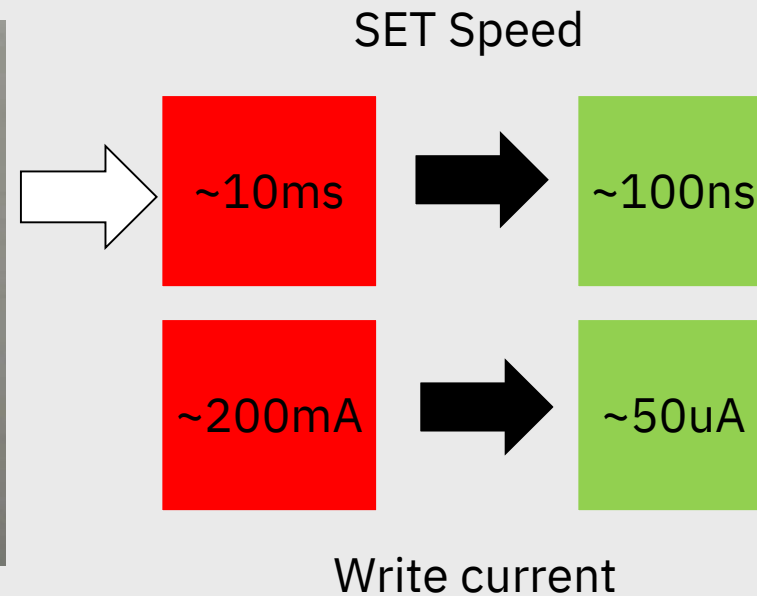
SET: 5mA, ~25V, 10ms

Read: 2.5mA, <5V

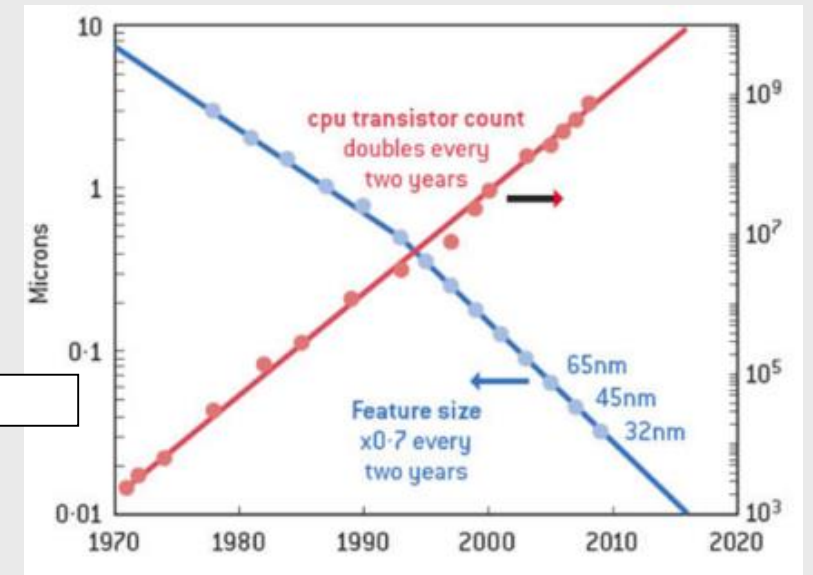


# A brief history of phase change memory

## Commercial success of optical recording



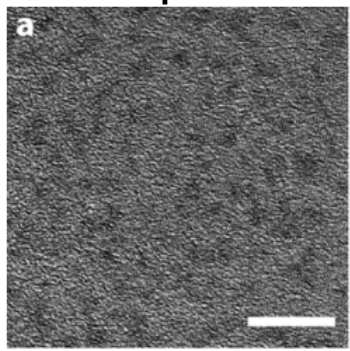
## Advances in semiconductor manufacturing



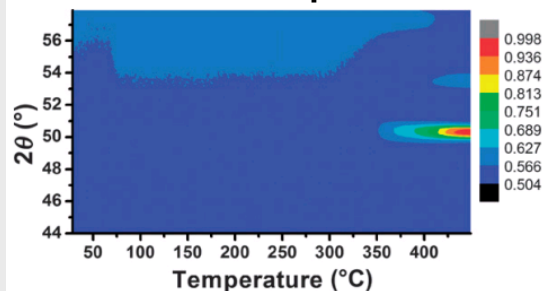
# Scalability

IBM ARC/YKT & Stanford  
(2007-2009)  
XRD studies of phase-change  
thin-films, nanodots,  
nanoparticles (1.8nm Ø)

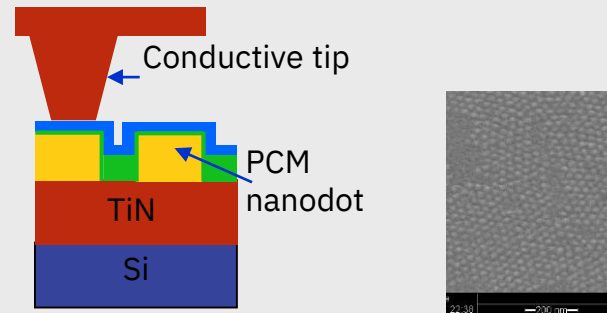
**TEM picture of 1.8nm  
nanoparticles**



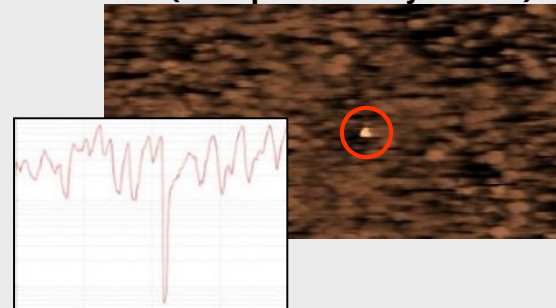
**In-situ XRD pattern**



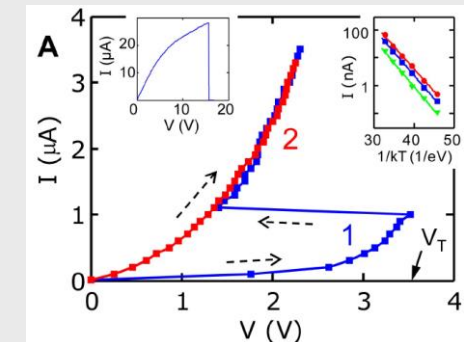
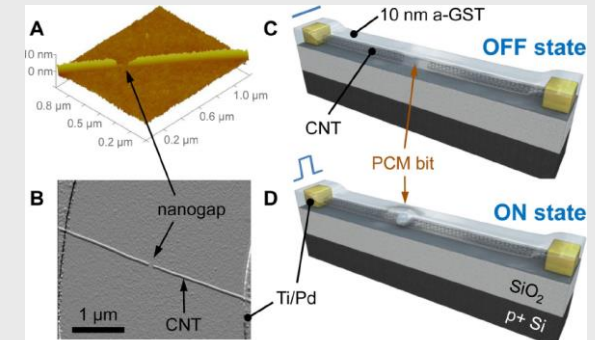
IBM-Zurich/Stanford (2009)  
Joule-heating induced  
switching of single phase-  
change nanodots (~15nm Ø)



**Switching of a single nanodot  
(amorphous → crystalline)**



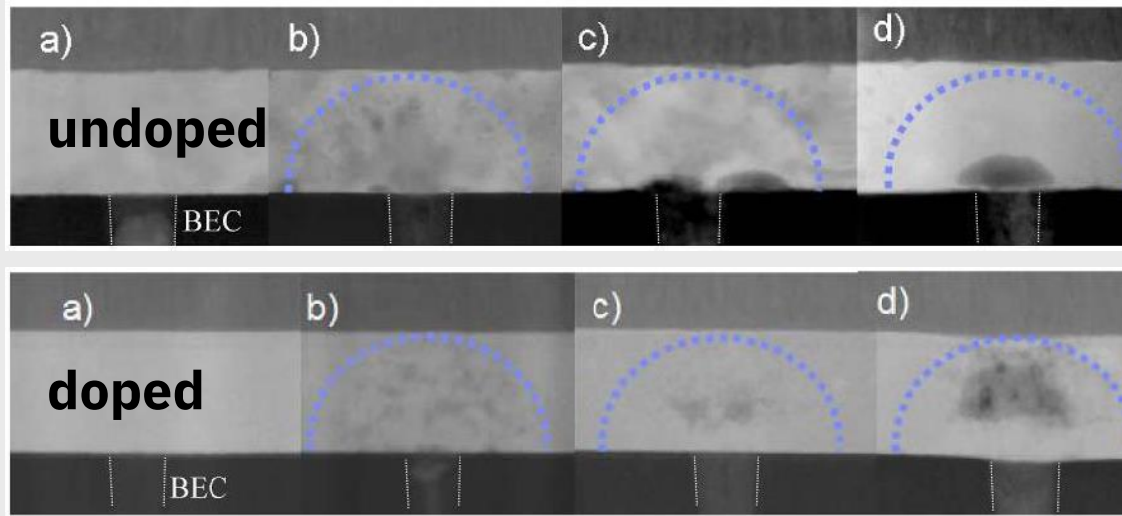
UIUC, Xiong et al., Science (2011)  
Phase change material deposited  
between a carbon nanotube which  
has been ruptured  
Nanotube diameter is ~2-3 nm



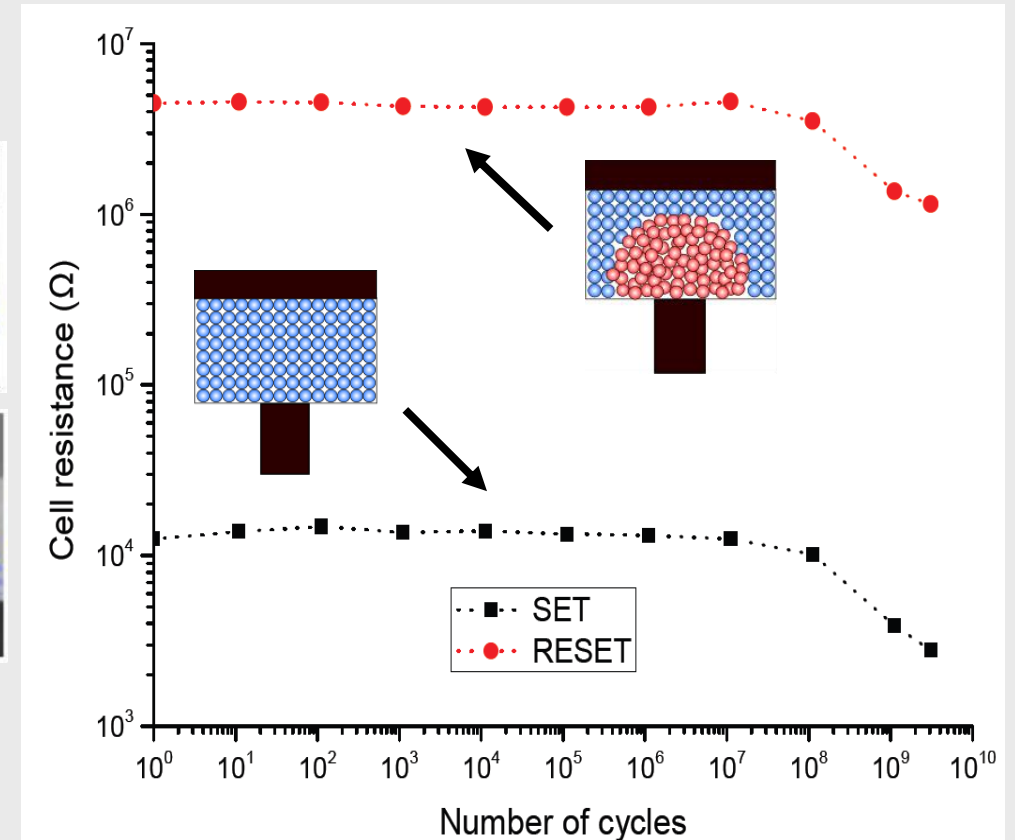
A phase change device can scale to a few nanometers

# Cycling endurance

## PCM cells with doped and undoped materials



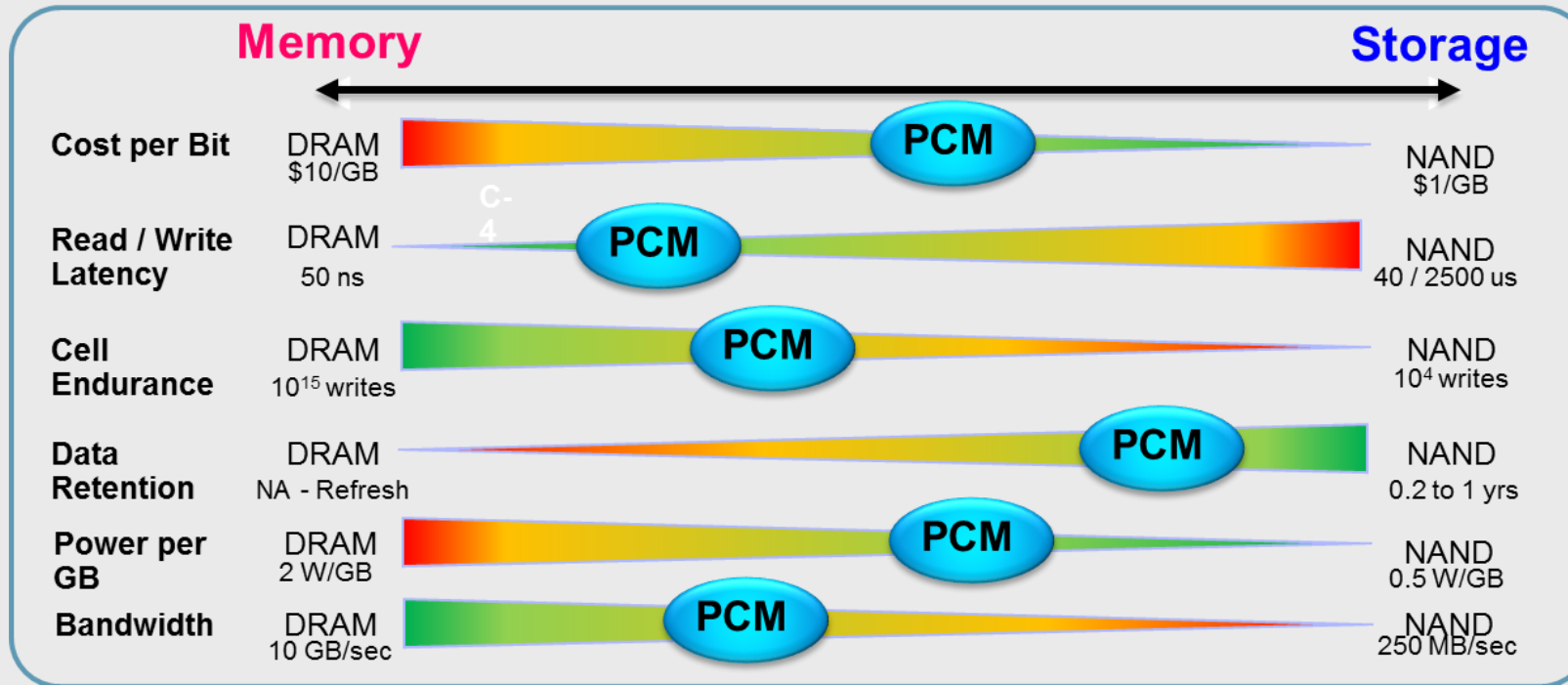
*Chen et al., IMW, 2009*



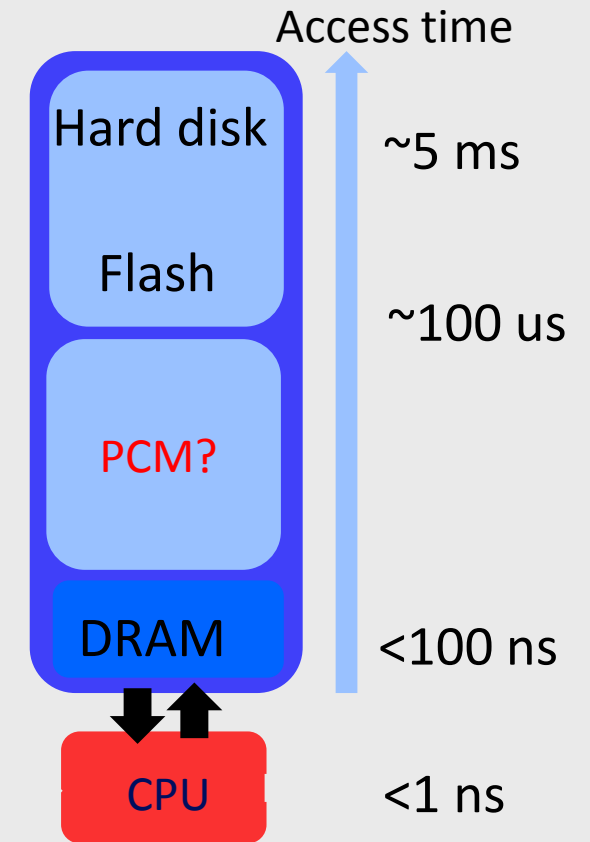
- Void formations and elemental segregation limit cycling endurance
- Doping/alloying the phase change material has shown to improve the endurance
- Cycling endurance of  $> 10^9$  feasible



# PCM as storage class memory

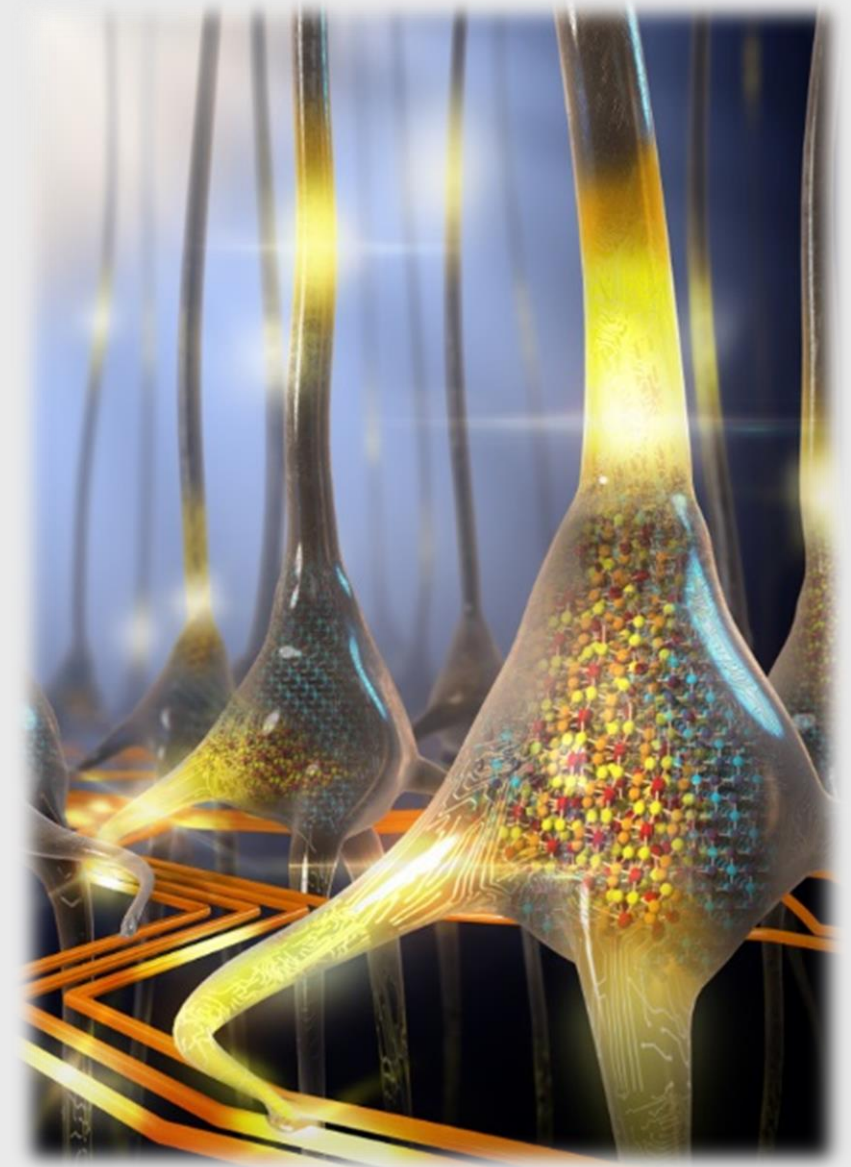


- Latency: much faster than FLASH (100's of ns vs. 100's of us)
- Write endurance: 1,000 x FLASH
- Nonvolatile, true random access capability, write in-place
- Very good scaling potential demonstrated (beyond 10nm node)
- Cost: between FLASH and DRAM (as technology matures)

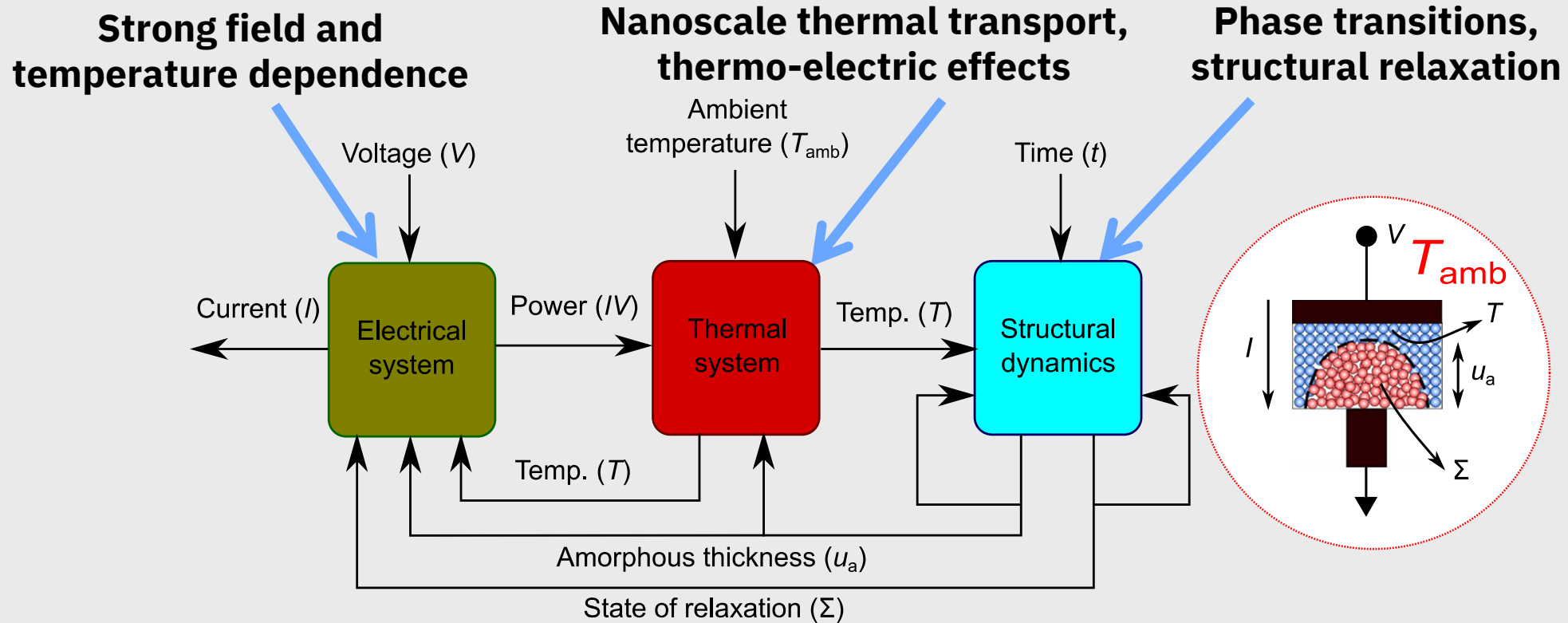


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# PCM device physics at a glance

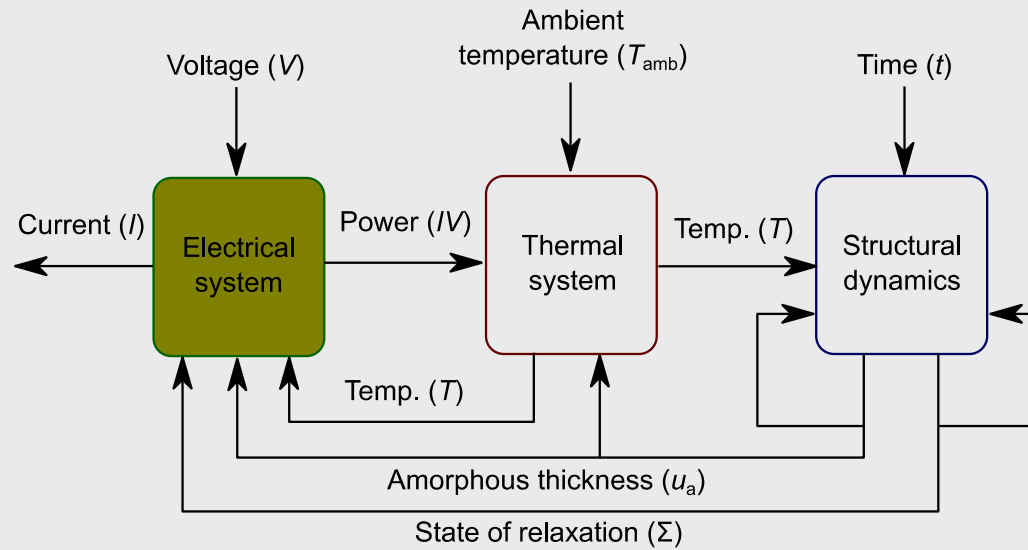


- Intricate feedback interconnection of electrical, thermal and structural dynamics
- Write operation: Alter the phase-configuration (Via Joule heating and structural dynamics)
- Read operation: Decipher the phase-configuration (Typically via reading the resistance at low field)

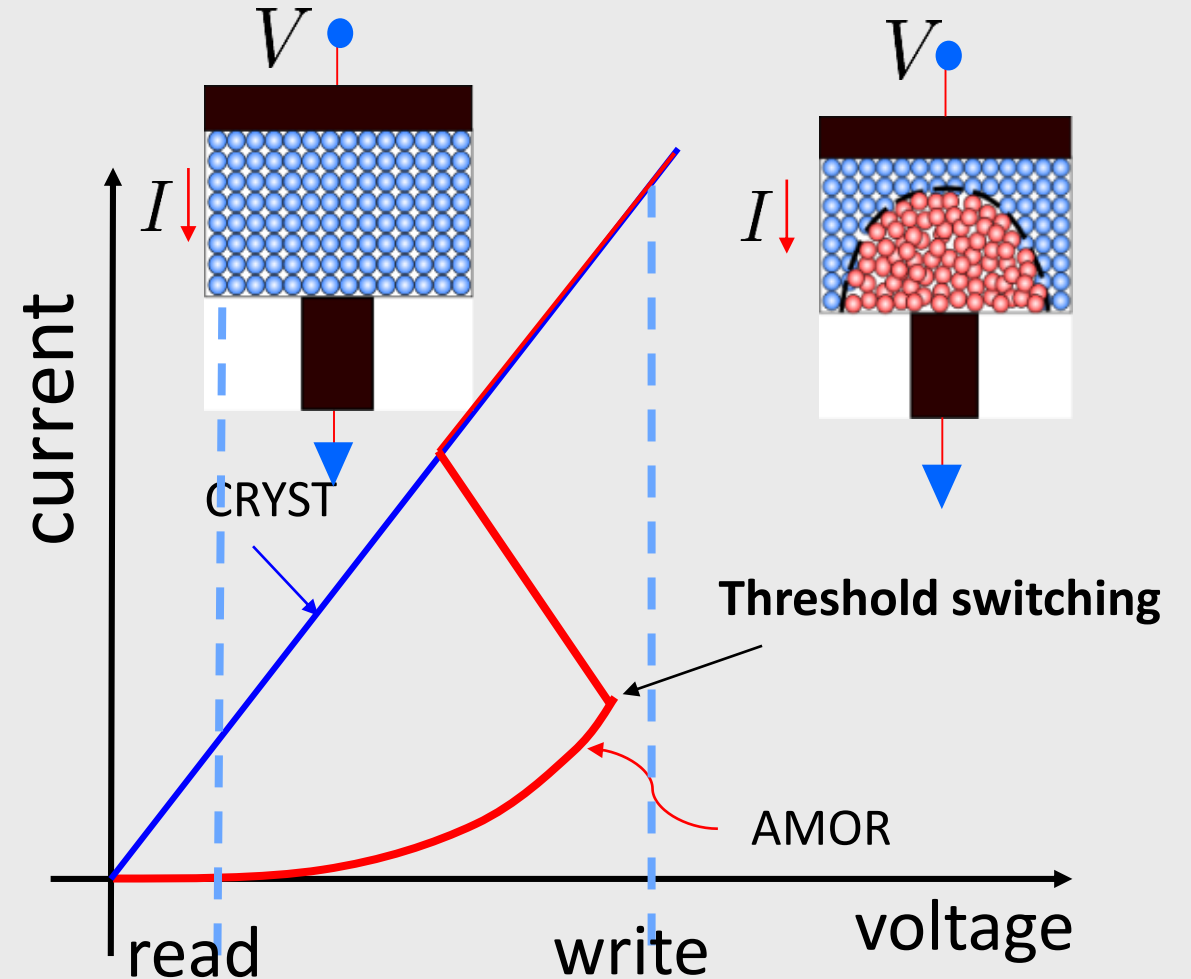
*Sebastian et al., Nature Comm., 2014, Le Gallo et al., New J. Phys., 2015, Le Gallo et al., J. Appl. Phys., 2016, Le Gallo et al., Adv. Electr., Mat., 2018*



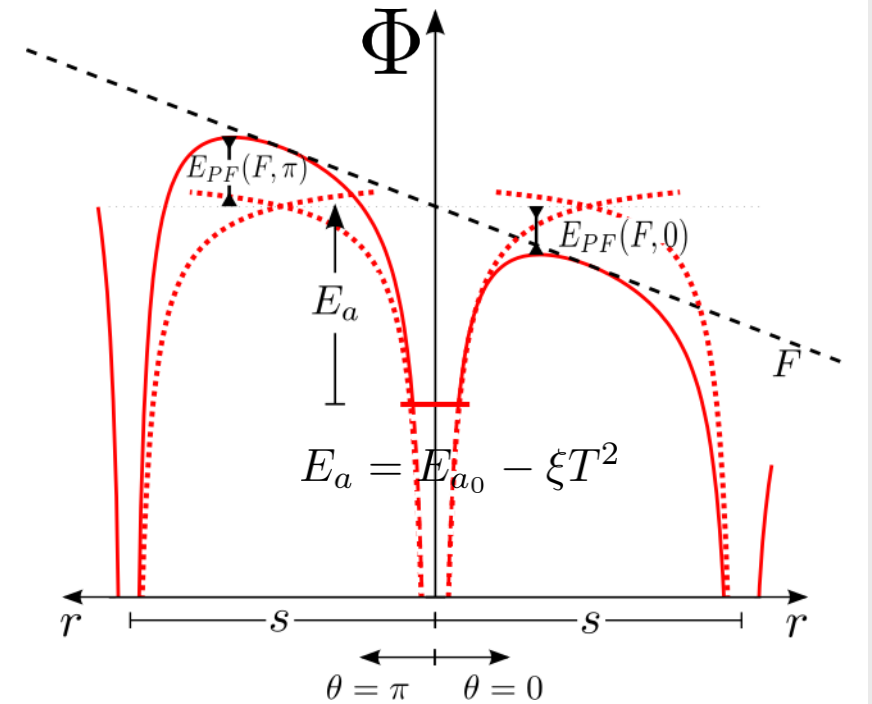
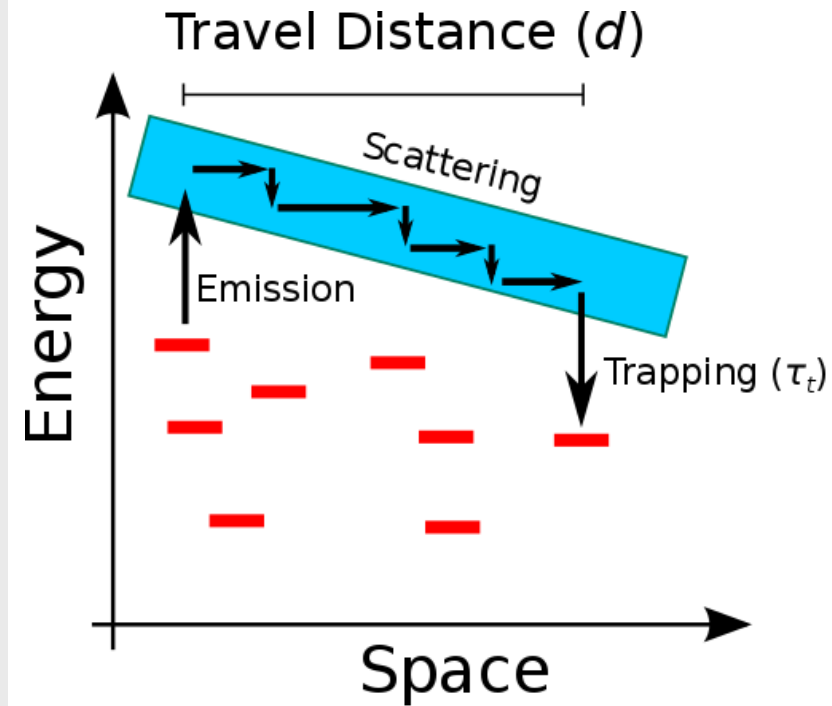
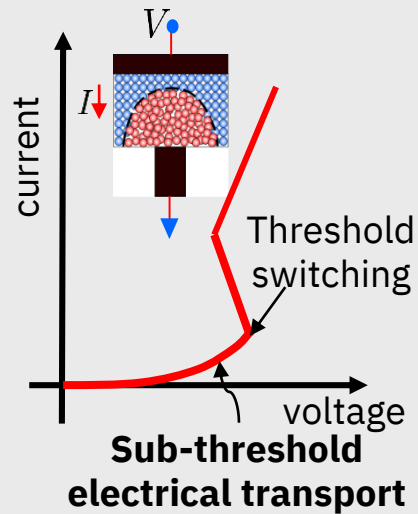
# PCM: The electrical system



- Crystalline phase exhibits near-Ohmic transport and has minimal temperature dependence
- Electrical transport in the amorphous phase exhibits strong temperature and field dependence
  - ✓ Sub-threshold region
  - ✓ Threshold switching



# Electrical transport model

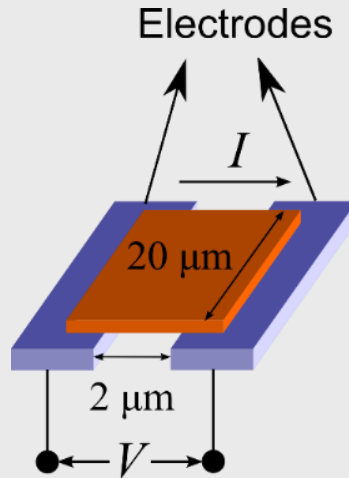


- Challenging to capture all the experimentally observed characteristics: 3 distinct regimes
  - ✓ Hill (1970), Hartke (1968), Ielmini (2007), Beneventi (2013)
- Well described by **trap limited band transport** together with **3D Poole-Frenkel emission** from a **two-center Coulomb potential**
- Two key parameters:
  - ✓ Activation energy for carrier emission:  $E_a$
  - ✓ The distance between the two defect centers:  $s$  (mean inter-trap distance)

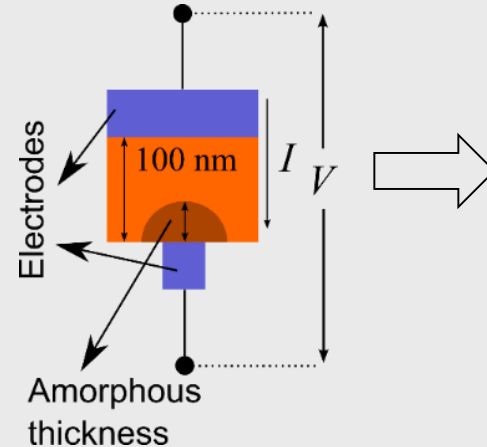
*Le Gallo et al., "Subthreshold electrical transport in amorphous phase-change Materials", New J. Phys., 2015*

# Model validation

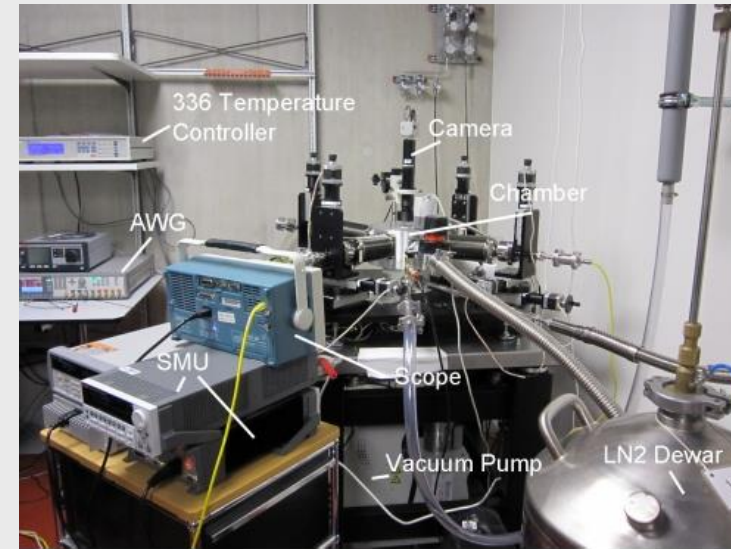
**Line Cell  
(as-deposited)**



**PCM Cell  
(melt-quenched)**



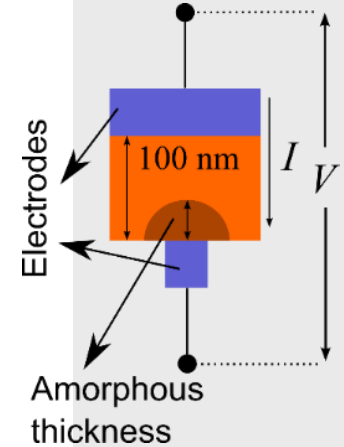
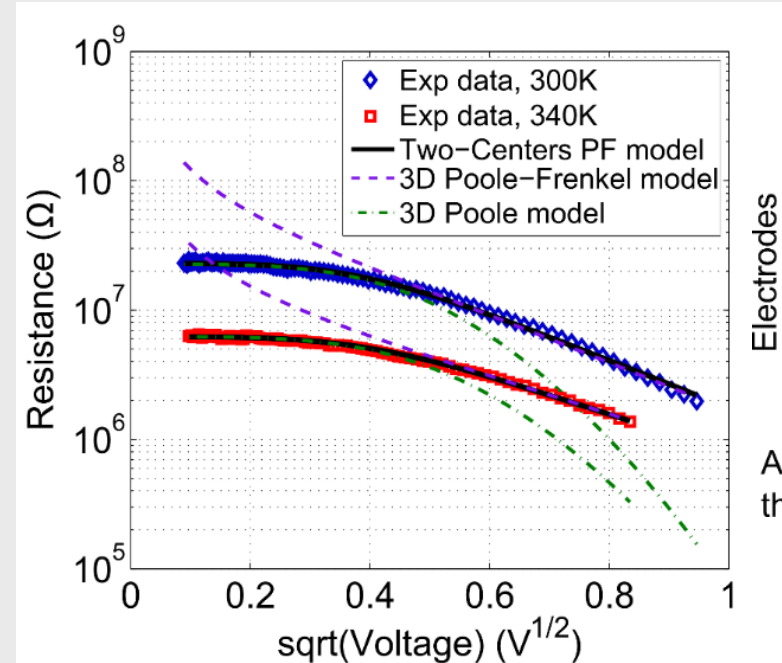
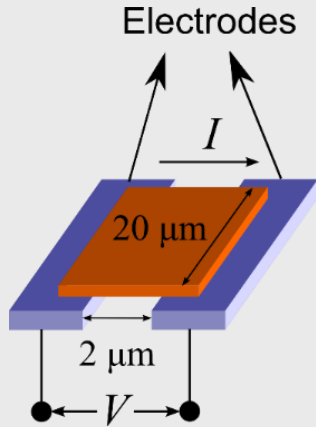
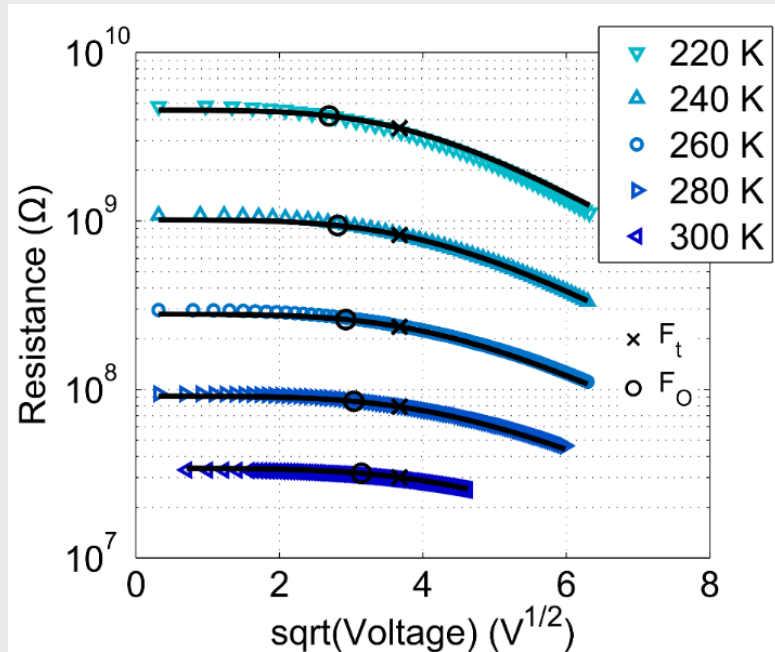
**Cryogenic probe station**



- IV curves measured on:
  - ✓ Line cells (uniform field distribution, large device of micrometer-scale)
  - ✓ PCM cells (non-uniform field distribution, small device of nanometer-scale)
- Cryogenic probe station used to obtain the IV characteristics over a wide temperature range



# Model validation



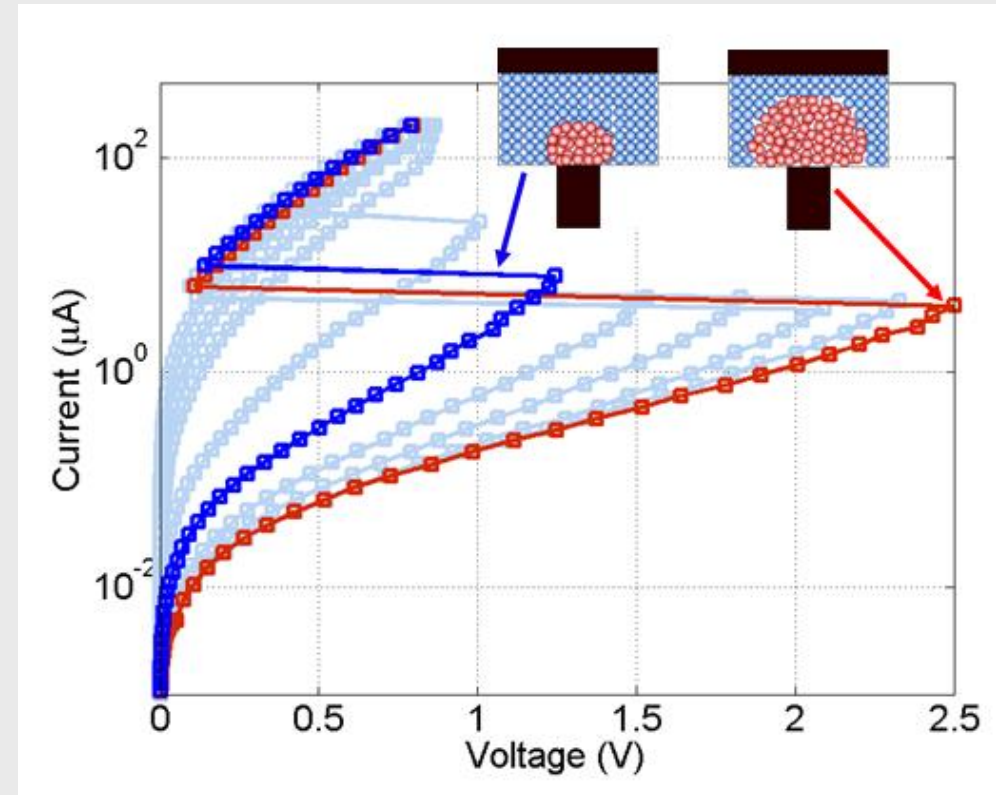
- $\epsilon_r = 13$  (FTIR)
- Electrode distance =  $2 \mu\text{m}$
- $E_a$  from  $R(T)$  data
- $s = 8.1 \text{ nm}$  (room temperature fit)

- $\epsilon_r = 13$  (FTIR)
- $U_{\text{Aeff}} = 15 \text{ nm}$
- $E_a$  from  $R(T)$  data
- $s = 8.1 \text{ nm}$  (room temperature fit)

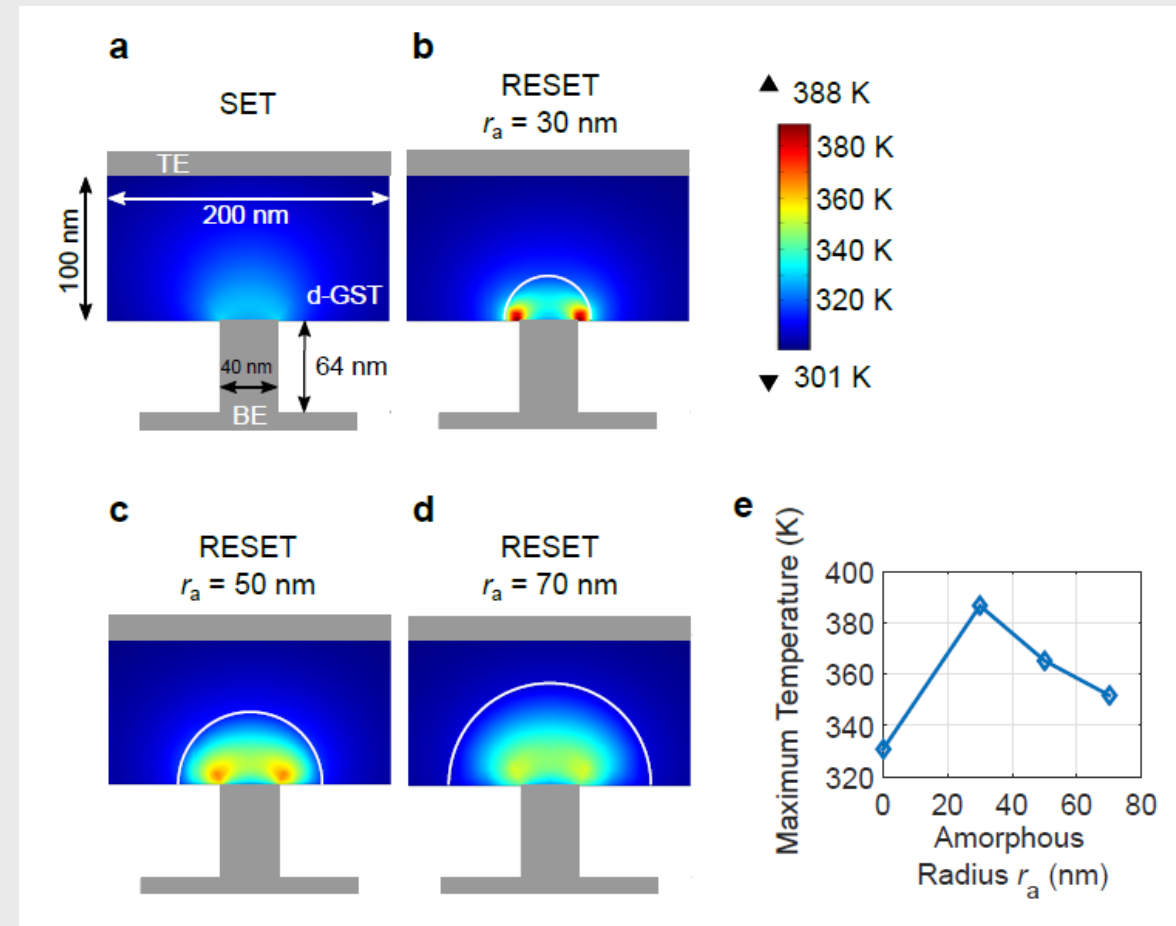
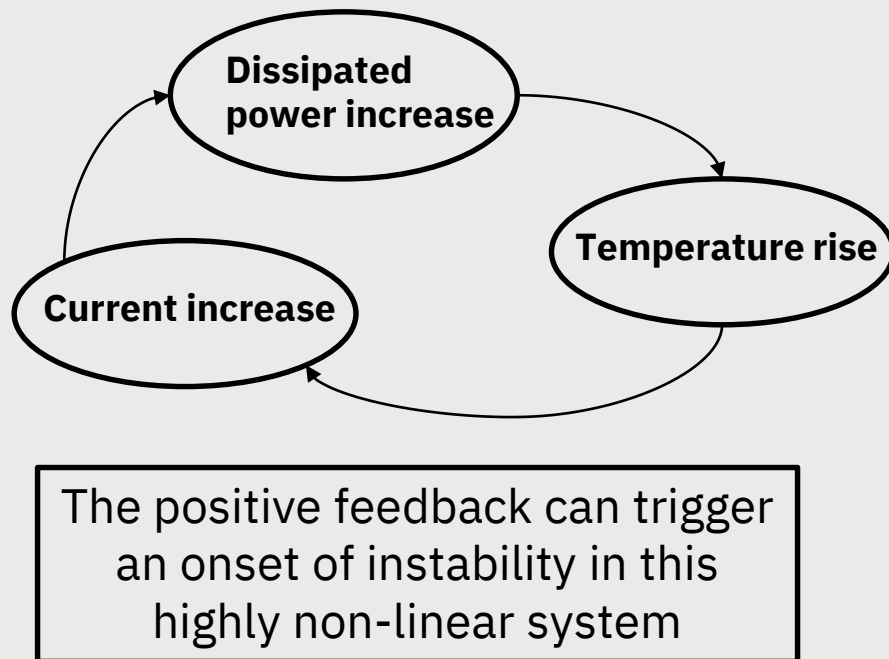
- Model captures the experimental data from 300K down to 220K without fitting across temperature
- Standard Poole and Poole-Frenkel models fail to capture the IV characteristics on the whole voltage range
- Full model captures the experimental data with similar physical parameters as for line cells

# Threshold switching

- Threshold switching enables dissipation of substantial electric power at relatively low voltages → A key enabling property
- Threshold switching voltage depends on the thickness of the amorphous region and the activation energy for electrical transport ( $E_a$ )
- After TS, the high field “ON” resistance similar for all states
- Details of the dynamics still a matter of ongoing research (after 50 years!)
  - ✓ **Thermal models** (Eaton, Boer, Tsendin etc.) vs **Purely electronic models** (Mott, Henisch, Adler, Pirovano, Lacaita etc.)
  - ✓ Others: Tunneling between trap states, energy gain via carrier temperature increase, field induced nucleation



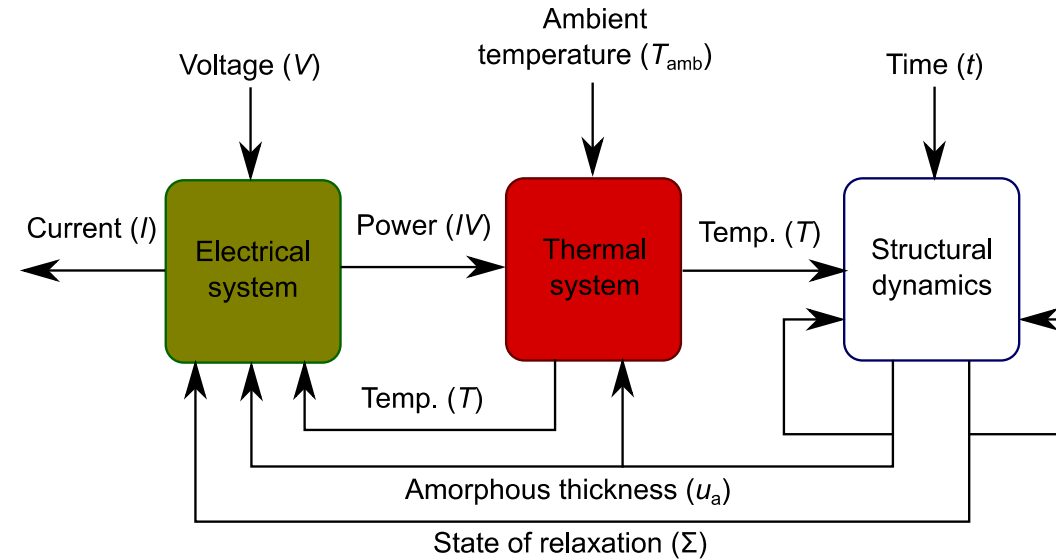
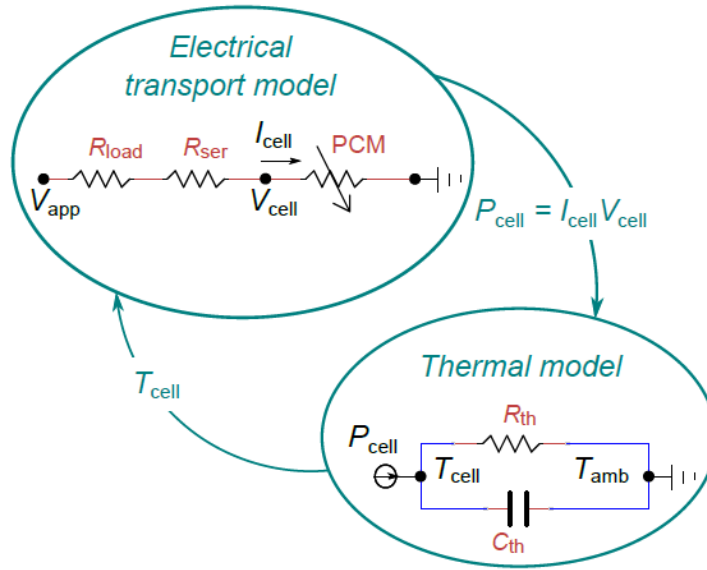
# Threshold switching in PCM devices: Thermally induced?



- High effective thermal resistance in nanoscale PCM devices ( $>1\text{K}/\mu\text{W}$ )
- Small thermal time constants ( $<10$  ns)



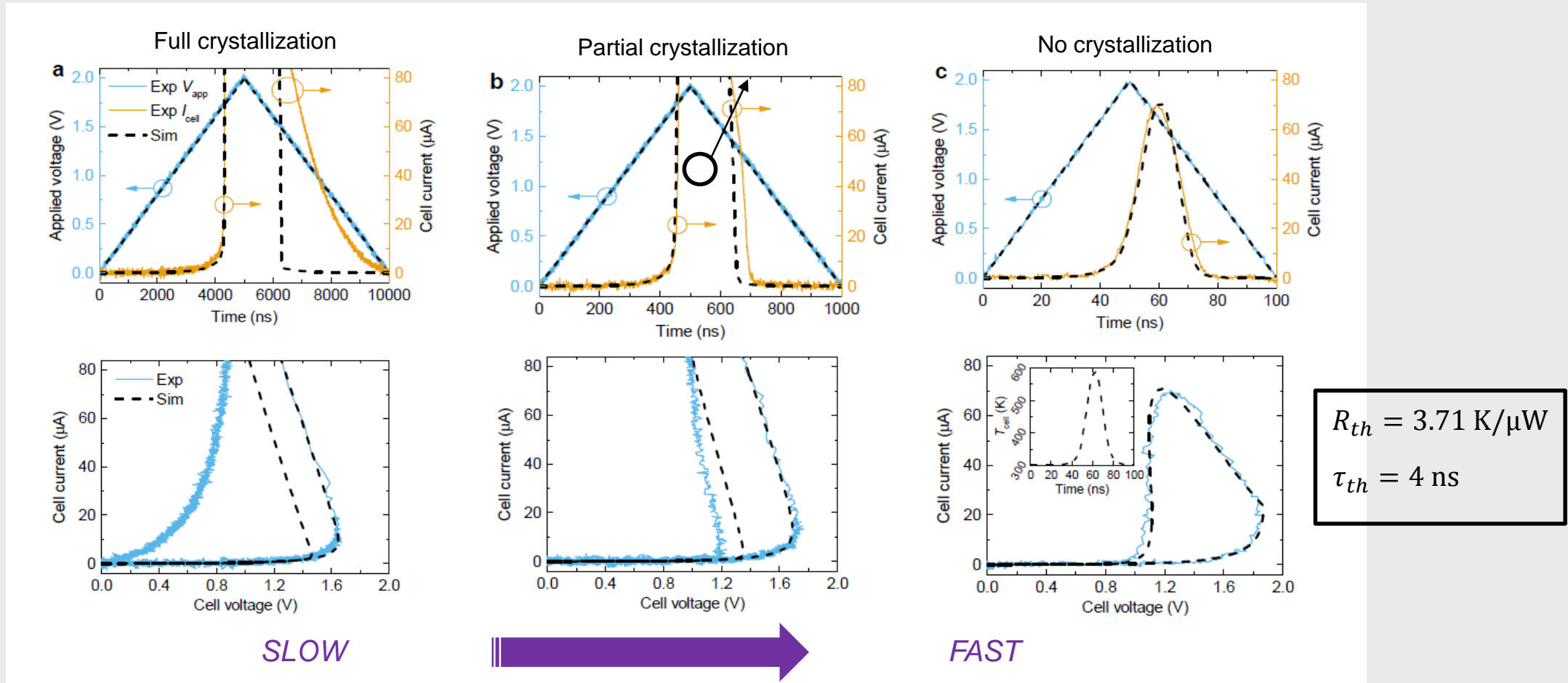
# Modeling



- Electrical transport model:
  - ✓ Field and temperature dependent subthreshold conduction model for PCM amorphous phase
  - ✓ Assume no influence of crystalline phase (e.g. no crystallization included)
- Thermal model:
  - ✓ Cell temperature rise from Joule heating coming from input power
  - ✓ Heat conduction described with effective thermal resistance and capacitance

***Le Gallo et. al, J. Appl. Phys., 2016***

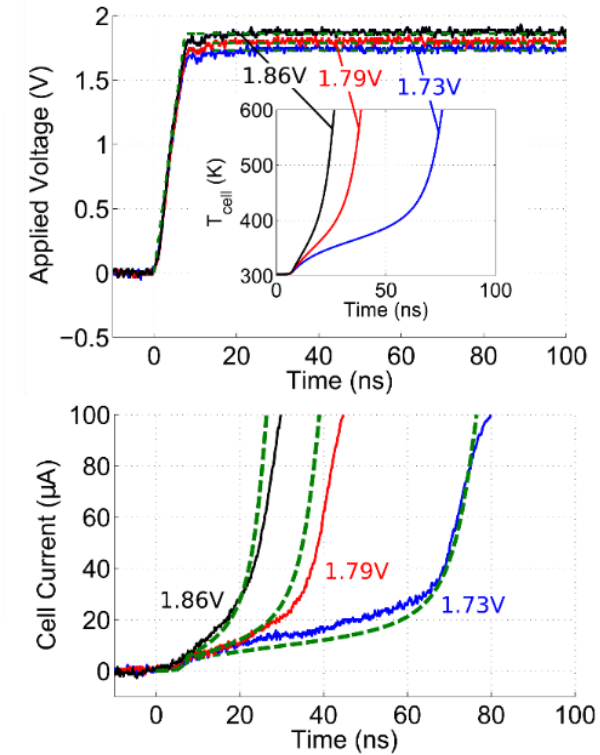
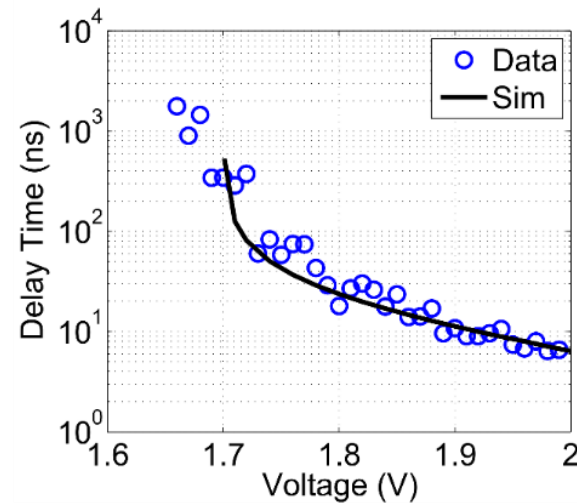
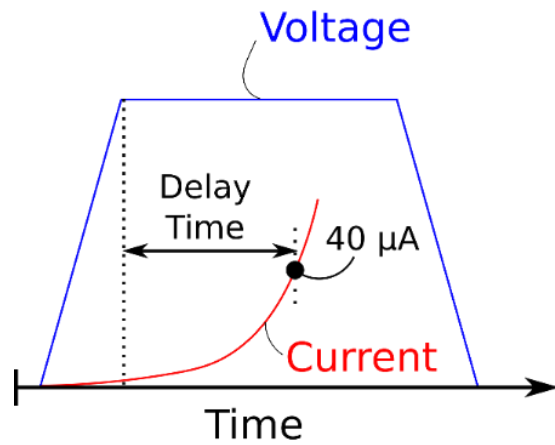
# Switching at Different Voltage Ramps



- Threshold switching dynamics are well captured by the thermally-assisted model
- Obtained values of  $R_{th}$  and  $\tau_{th}$  realistic for nanoscale PCM cells.

*Le Gallo et. al, J. Appl. Phys., 2016*

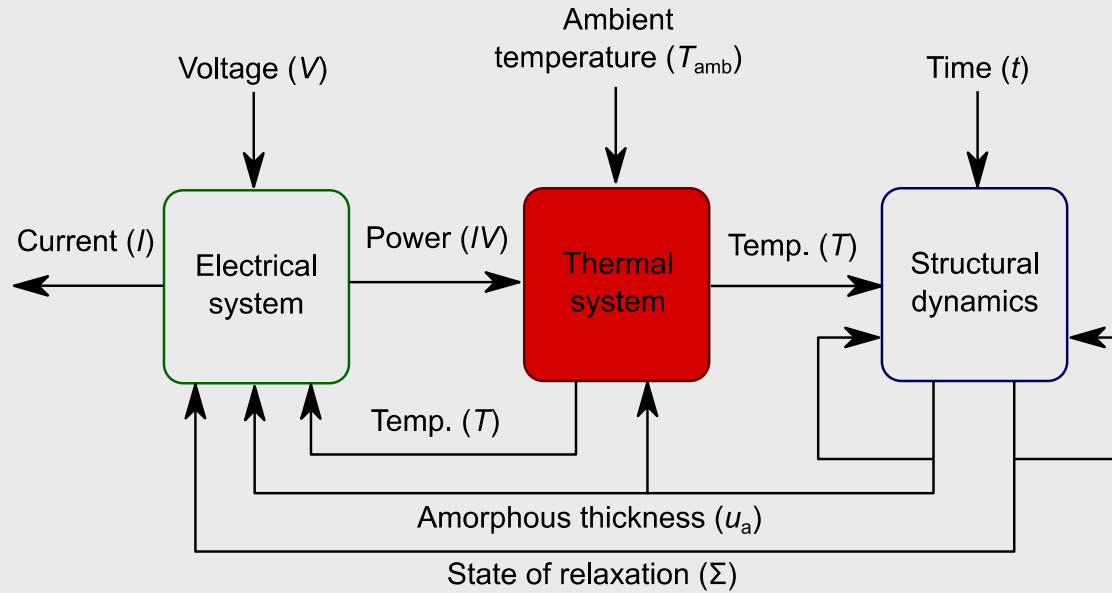
# Delay Time Switching



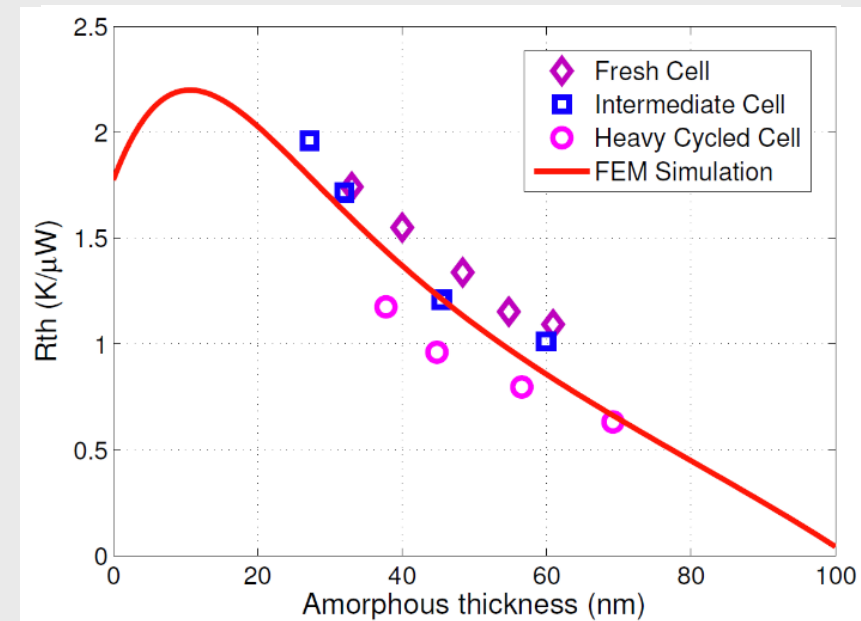
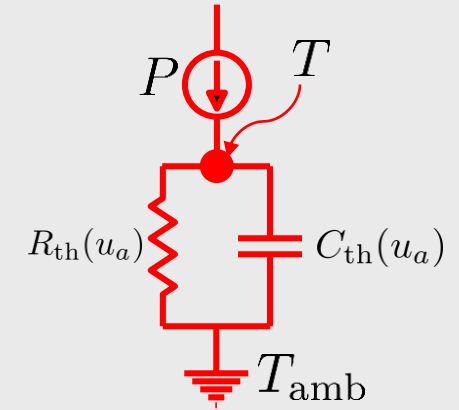
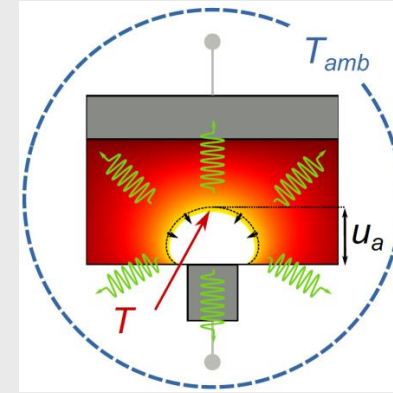
- Delay-time switching with voltage box pulse
- Short delay times ( $<1000\ \text{ns}$ ) well captured by thermal feedback model
- Slow current rise prior to switching explained by temperature buildup
- Manifestations of feedback instability in a highly non-linear system!

*Le Gallo et. al, J. Appl. Phys., 2016*

# PCM: The thermal system



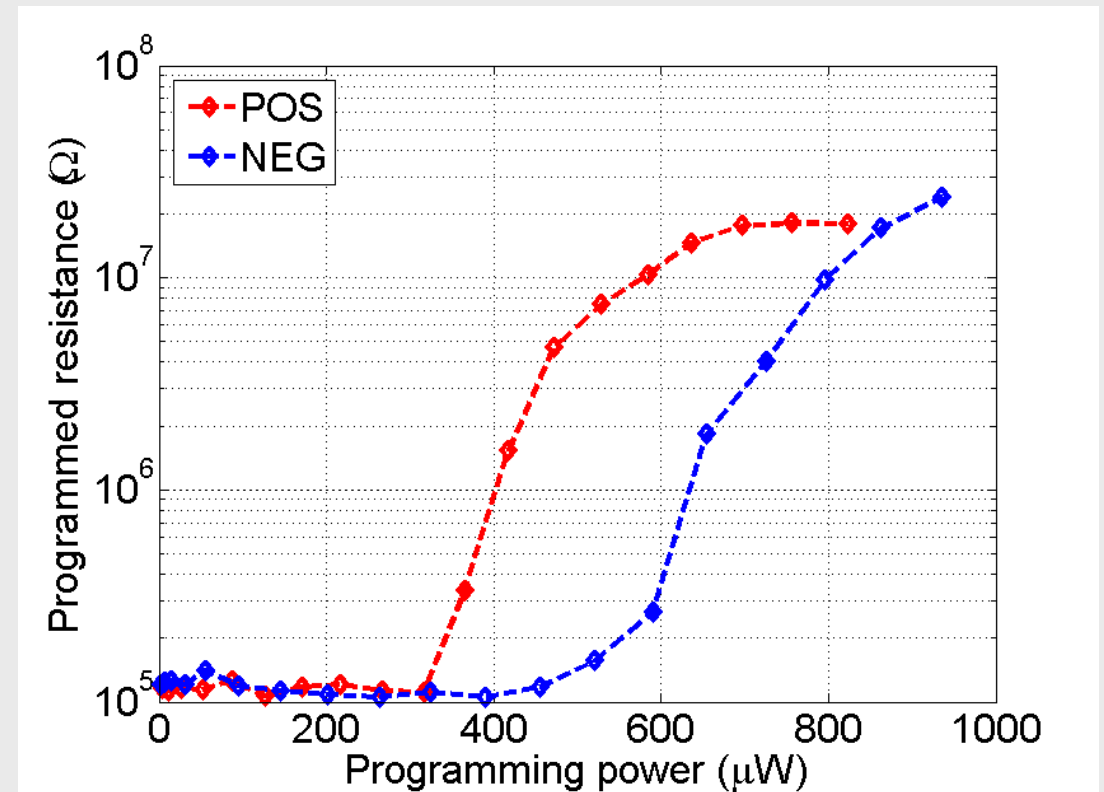
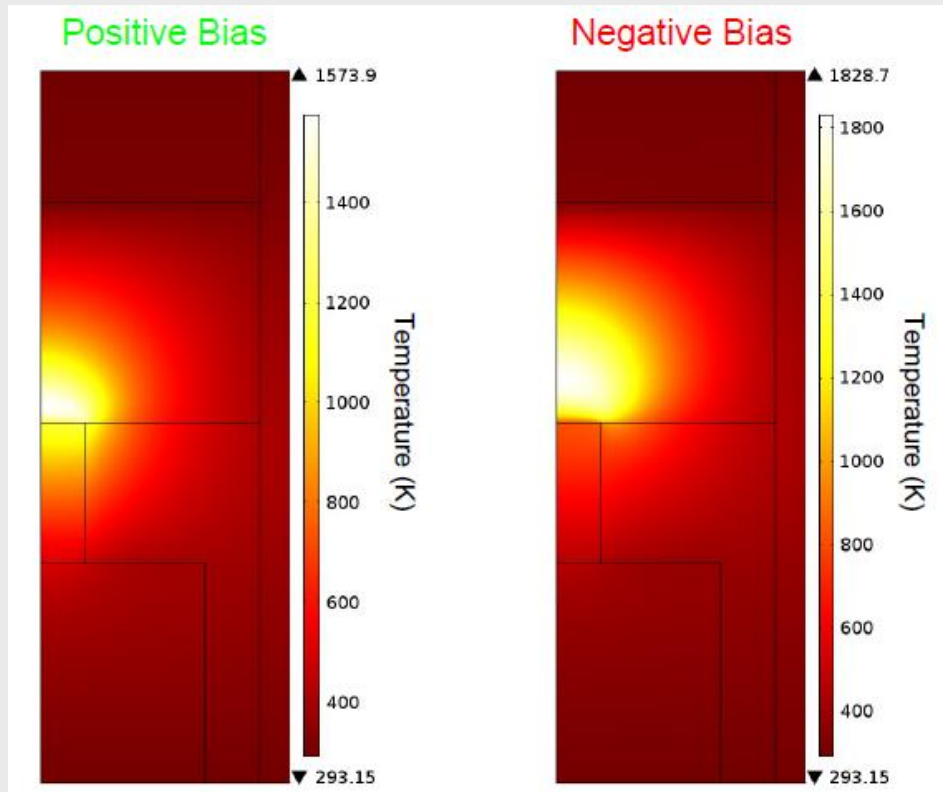
- The electrical power that is dissipated in the device results in a temperature rise within the device
- Temperature distribution with a “Hot spot” within the PCM segment
- Temperature distributions in a PCM device can be simulated and even experimentally verified!
- The “thermal” dynamics can be captured by an equivalent thermal resistance and capacitance



**Sebastian et. al, Nature Comm., 2014**



# Thermo-electric Effects

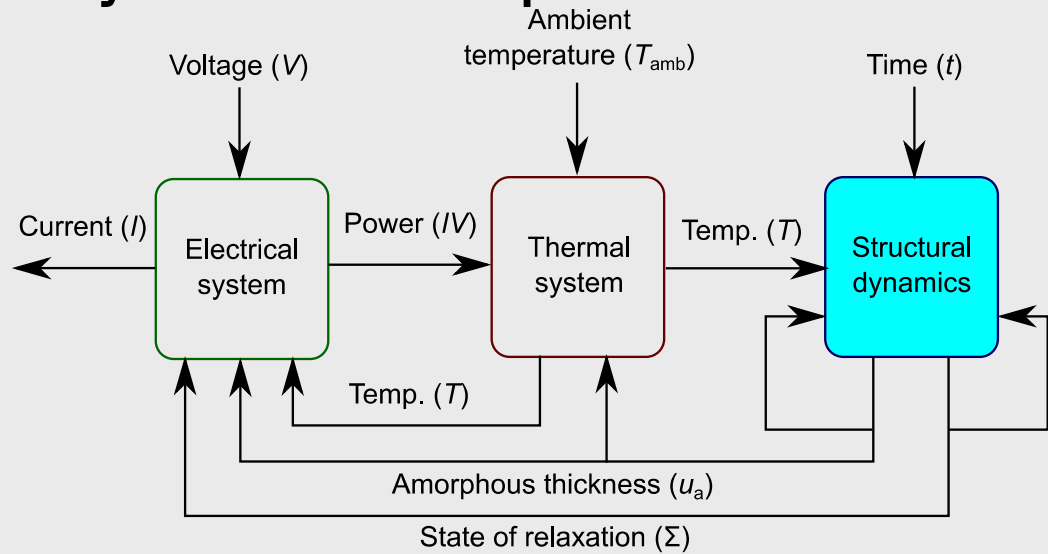


- The large temperature gradients  $\rightarrow$  thermo-electric effects such as Seebeck effect and the non-Joule heating terms such as Thomson, Peltier and Bridgeman cannot be neglected
- Polarity dependence in asymmetric devices

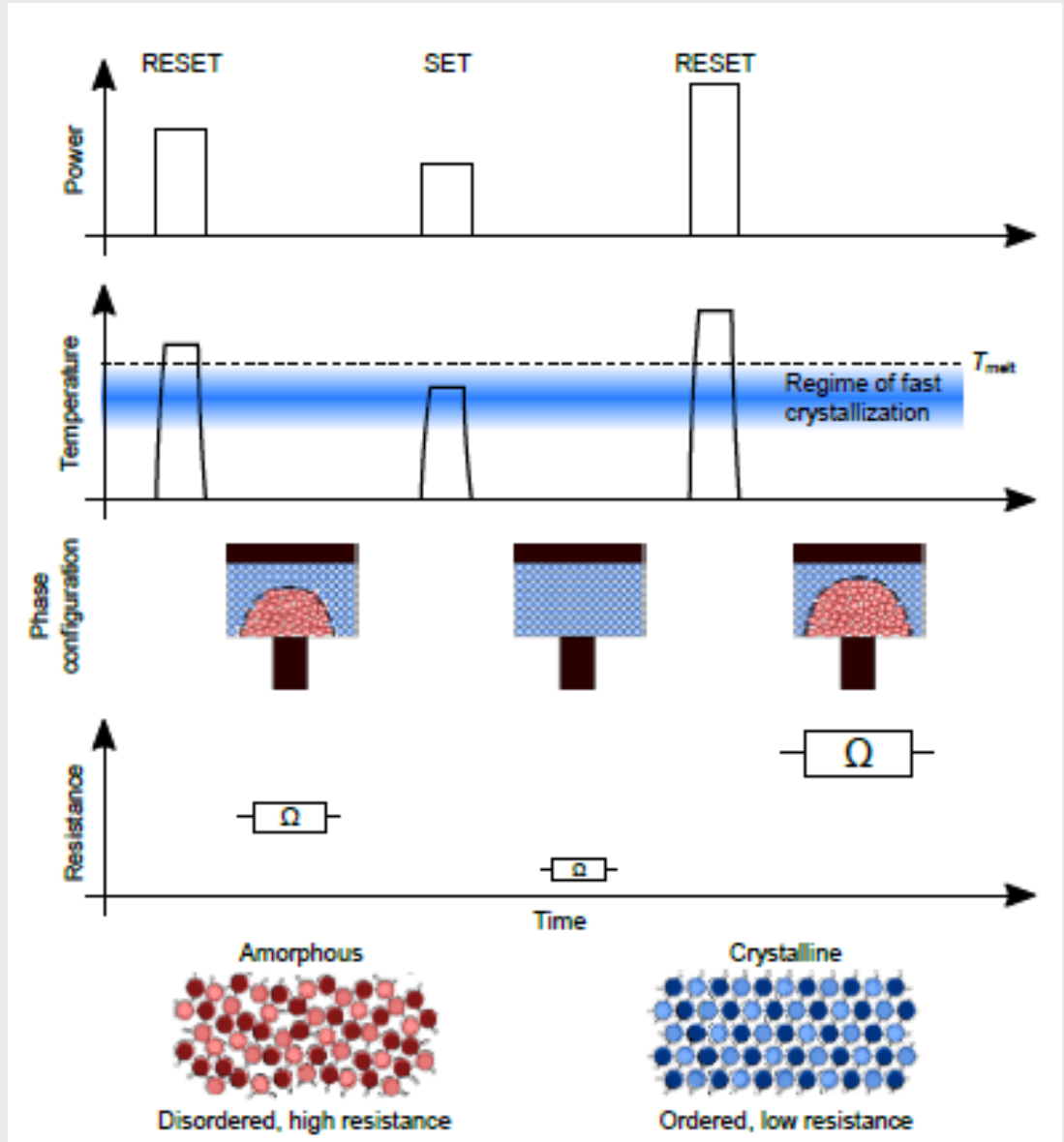
*Lee et al., Nanotechnology, 2012*

*Athmanathan et al., SISPAD, 2015*

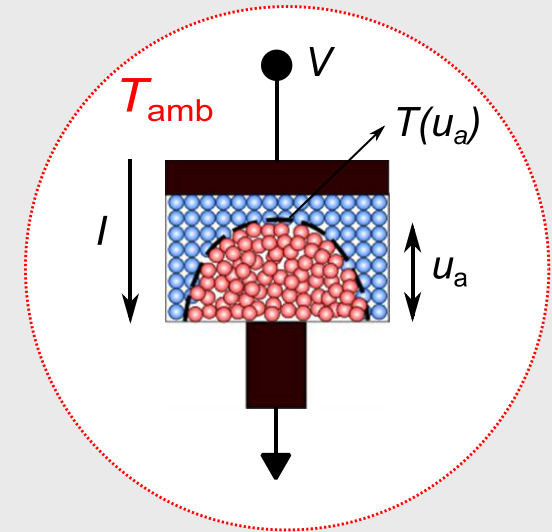
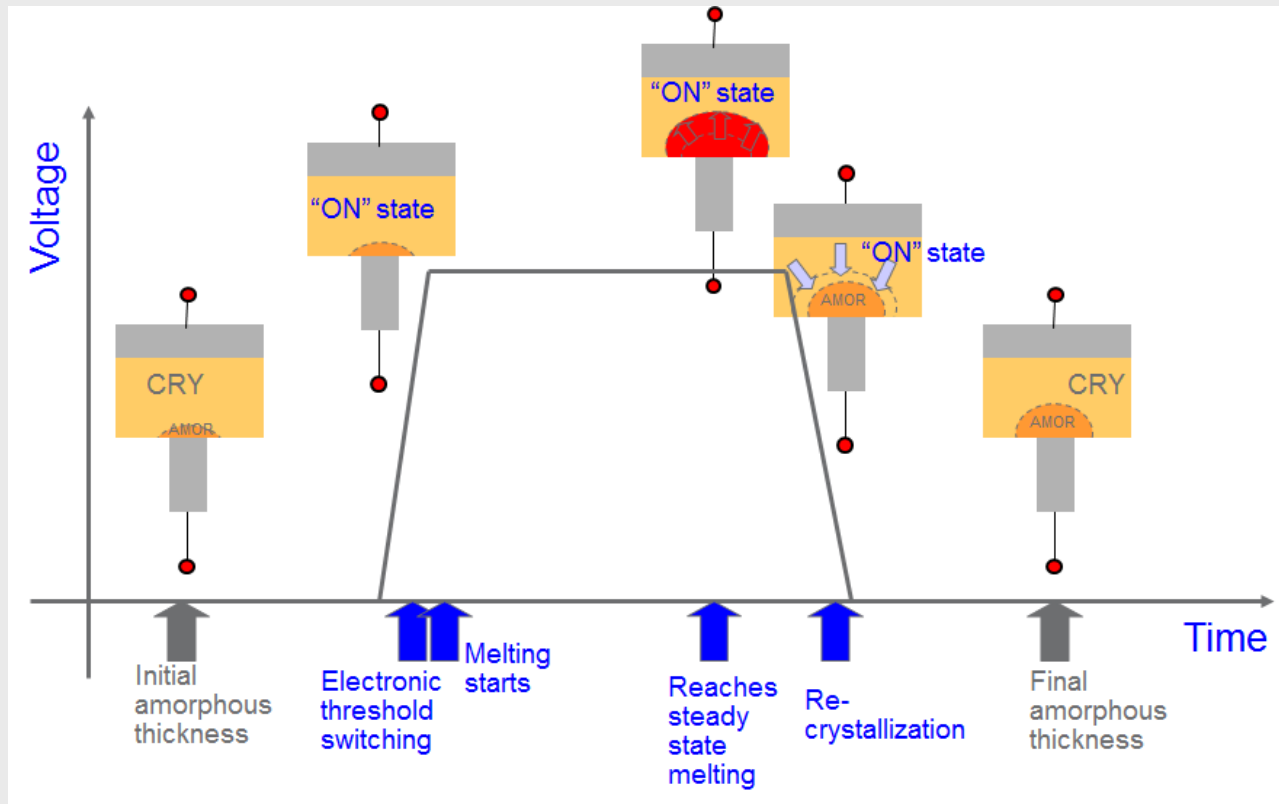
# Structural changes induced by the write process



- During the write process we induce structural changes to the phase-change material
- The structural dynamics is governed mostly by the temperature distribution within the device
- **Melting and quenching** to create disordered amorphous phase
- **Crystallization** to create ordered crystalline phase



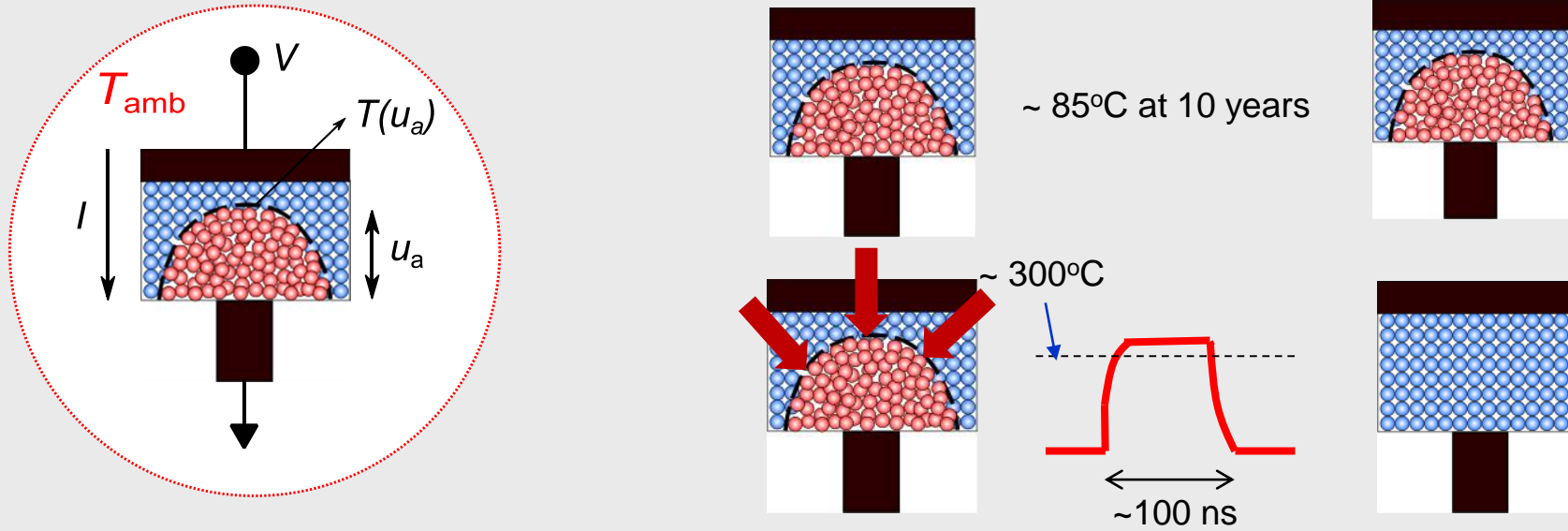
# PCM structural dynamics: “Amorphization” via Melt-quench



$$\frac{du_a}{dt} = \frac{1}{2\pi u_a^2} \left[ \frac{T(u_a) - T_m}{R_{th}(u_a) \Delta H} \right]$$

- Melting kinetics mostly governed by the thermal resistance ( $R_{th}$ ) and latent heat of fusion ( $\Delta H$ )
- $T_m$  is the melting temperature of the phase change material (approx. 890 K for  $\text{Ge}_2\text{Sb}_2\text{Te}_5$ )
- The melting process takes some time (in the order of tens of nanoseconds) and eventually  $T(u_a)$  equals  $T_m$  where a steady state will be achieved and no further melting is possible
- The molten material is cooled down abruptly to freeze the atomic structure into a disordered state

# PCM structural dynamics: Crystal growth

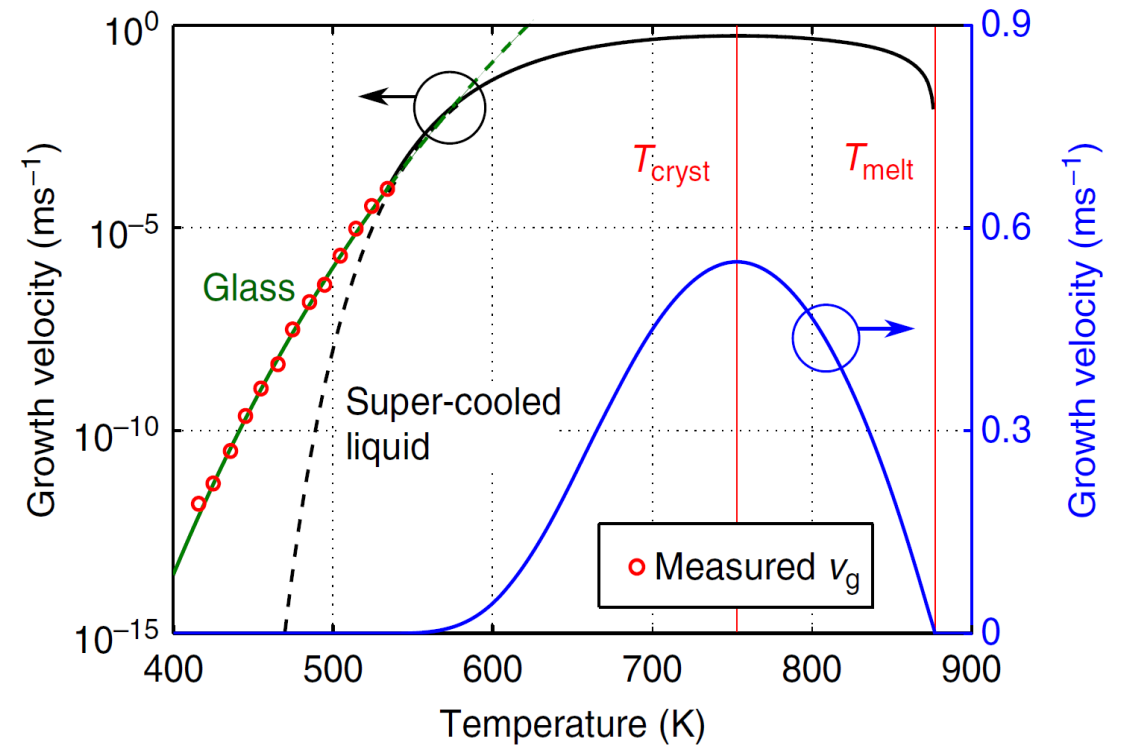


- Crystallization in nanoscale devices dominated by crystal growth
- Moreover, melt-quench phase has a significant number of built-in nuclei
- Crystal growth in phase change materials changes by 16 to 17 orders of magnitude as a function of temperature
- Single most important property that enables PCM



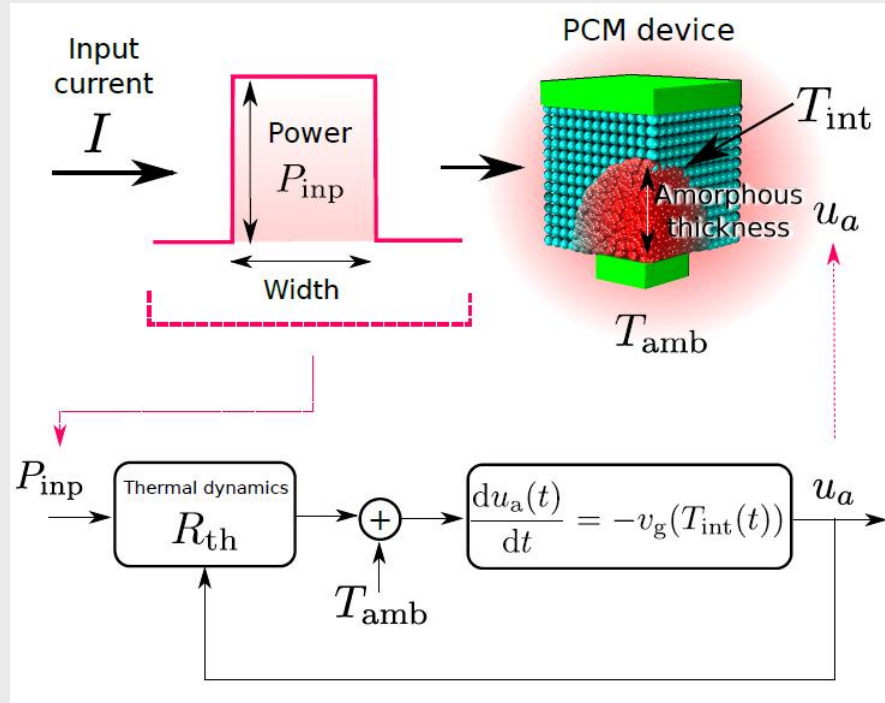
# Temperature dependence of crystal growth

- First experimental measurement of T dependence of growth velocity till melting temperature
- Measurements made directly inside the PCM device!
- Arrhenius-like temperature dependence spanning over 8 orders of magnitude
  - ✓ Deviates from the super-cooled liquid behavior potentially due to the ultra-fast quench rate!



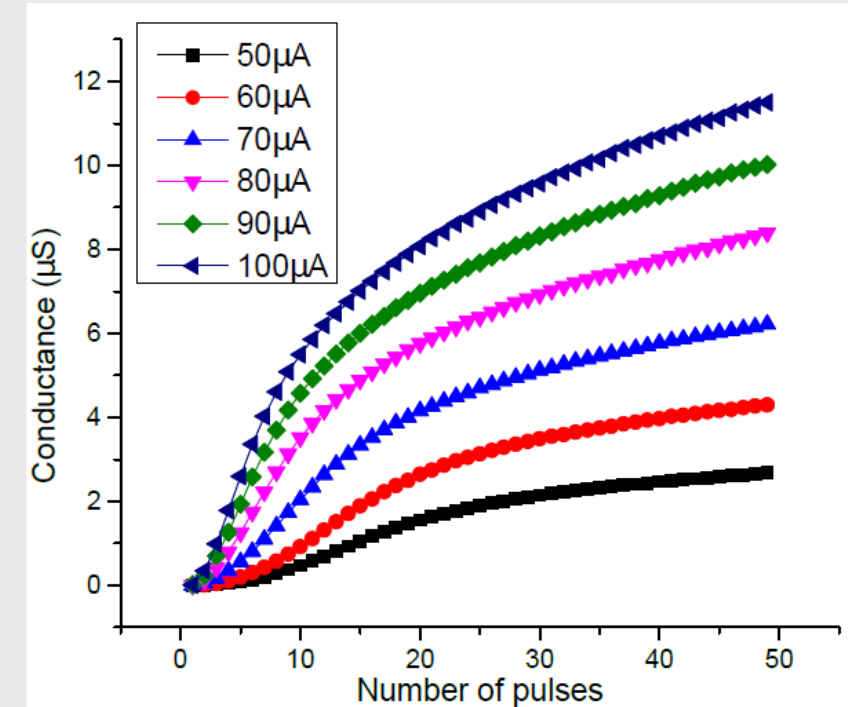
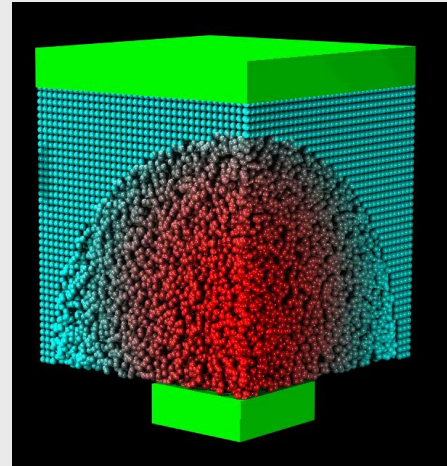
**Sebastian et. al, Nature Comm., 2014**

# Accumulative behavior



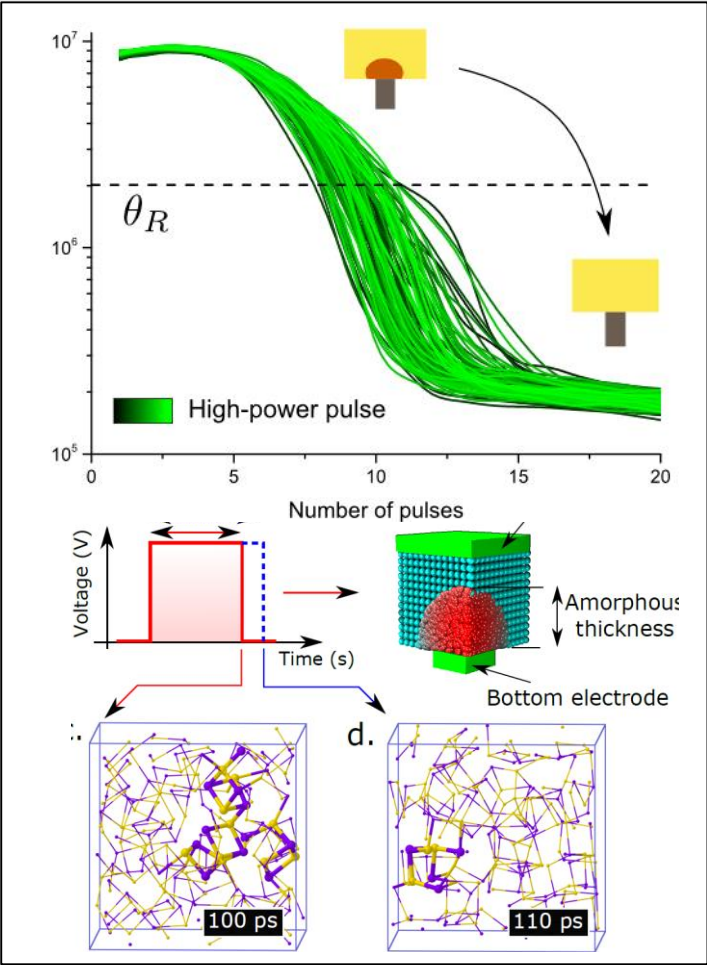
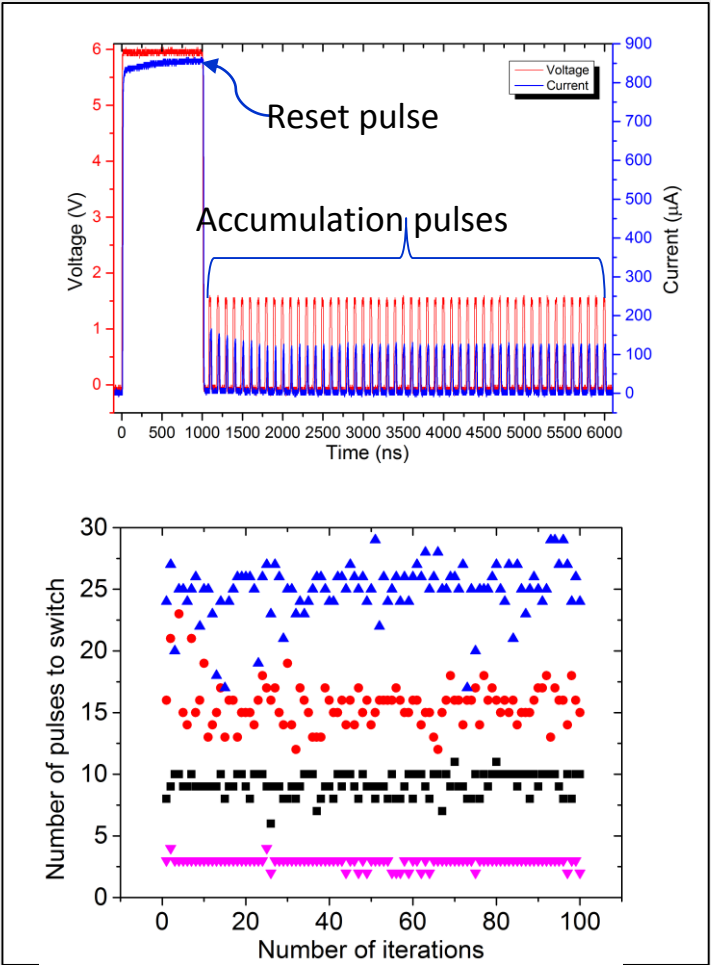
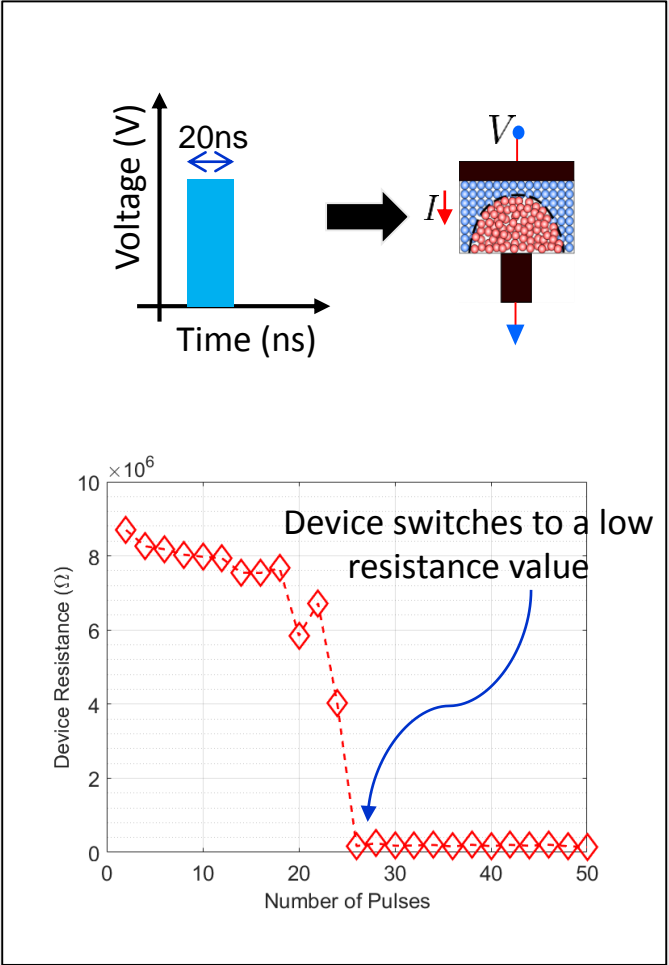
$$\frac{du_a(t)}{dt} = -v_g(T(t))$$

$$u_{a0} = u_a(0)$$

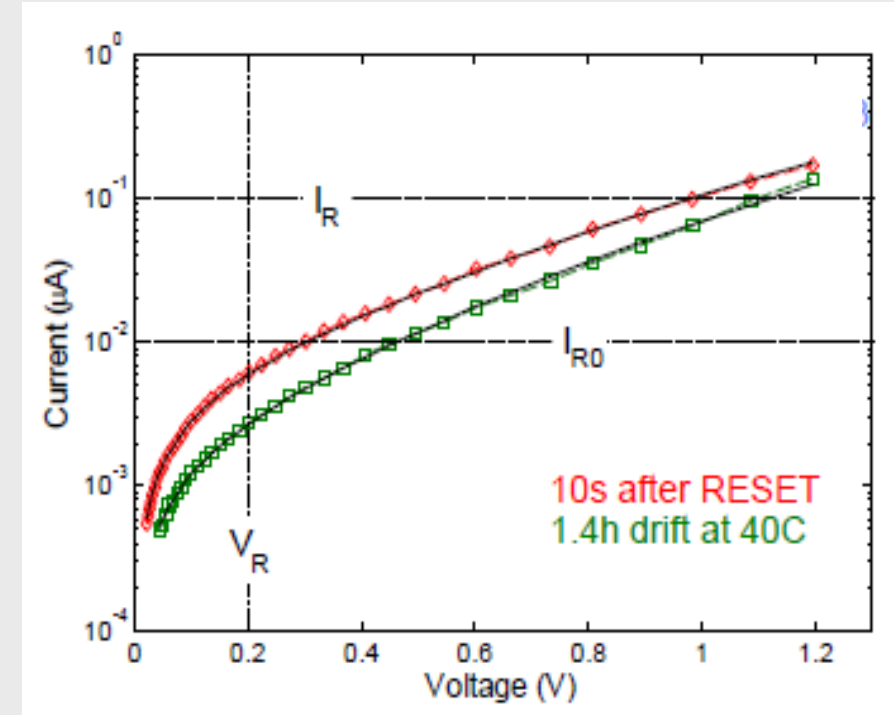
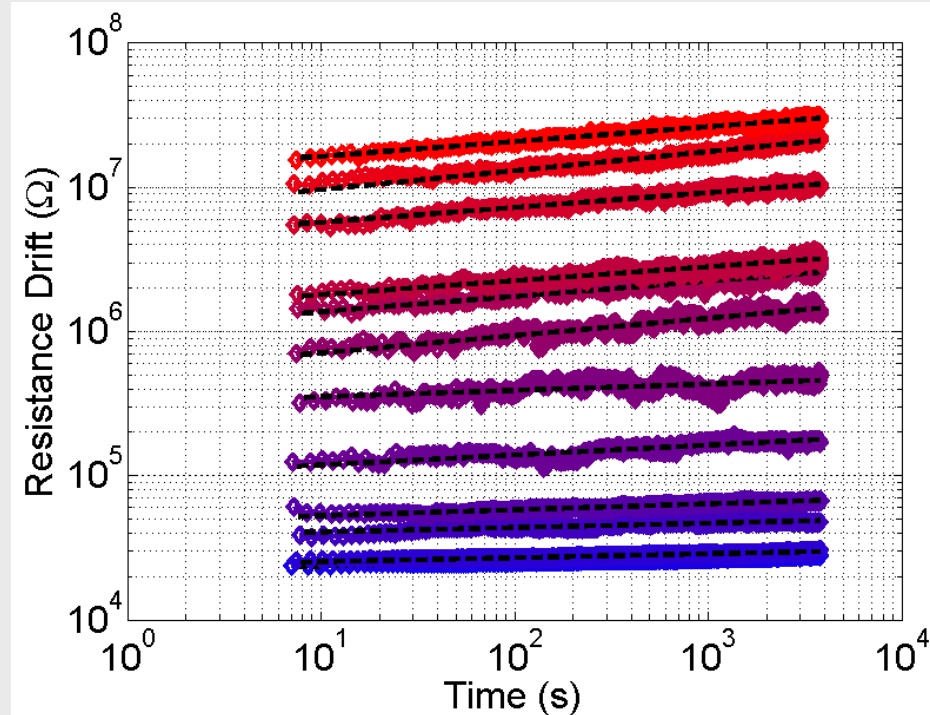


*Sebastian et. al, Nature Comm., 2014*

# Inherent randomness



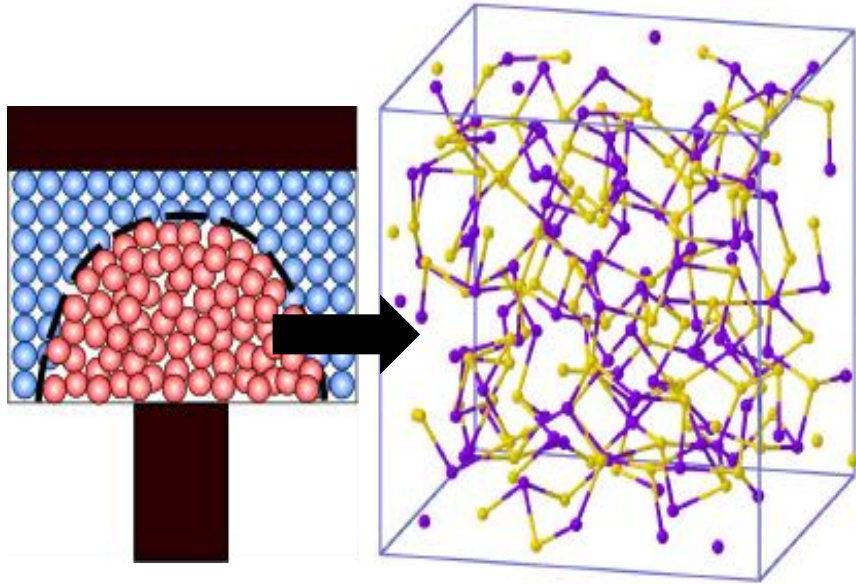
# “Drift” in PCM



*Ielmini et al., IEEE Trans. Elec. Dev., 2009*  
*Rizzi et al., Appl. Phys. Lett., 2011*  
*Fantini et al., Appl. Phys. Lett., 2012*

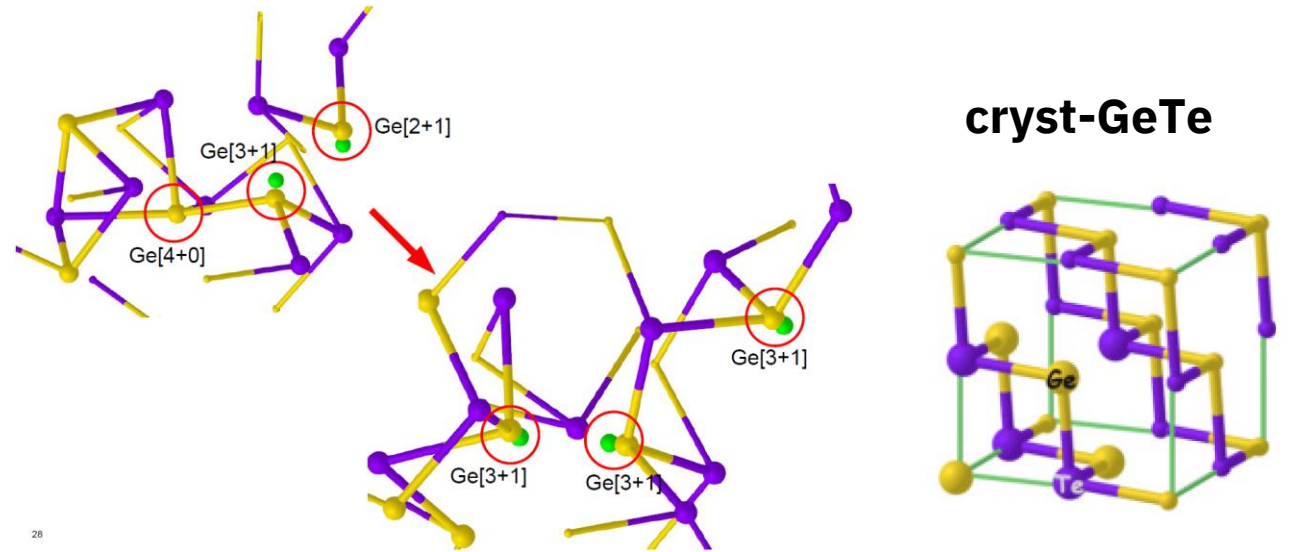


# The cause of drift: Spontaneous structural relaxation



Structural relaxation of amorphous phase  
towards an “ideal” glass state

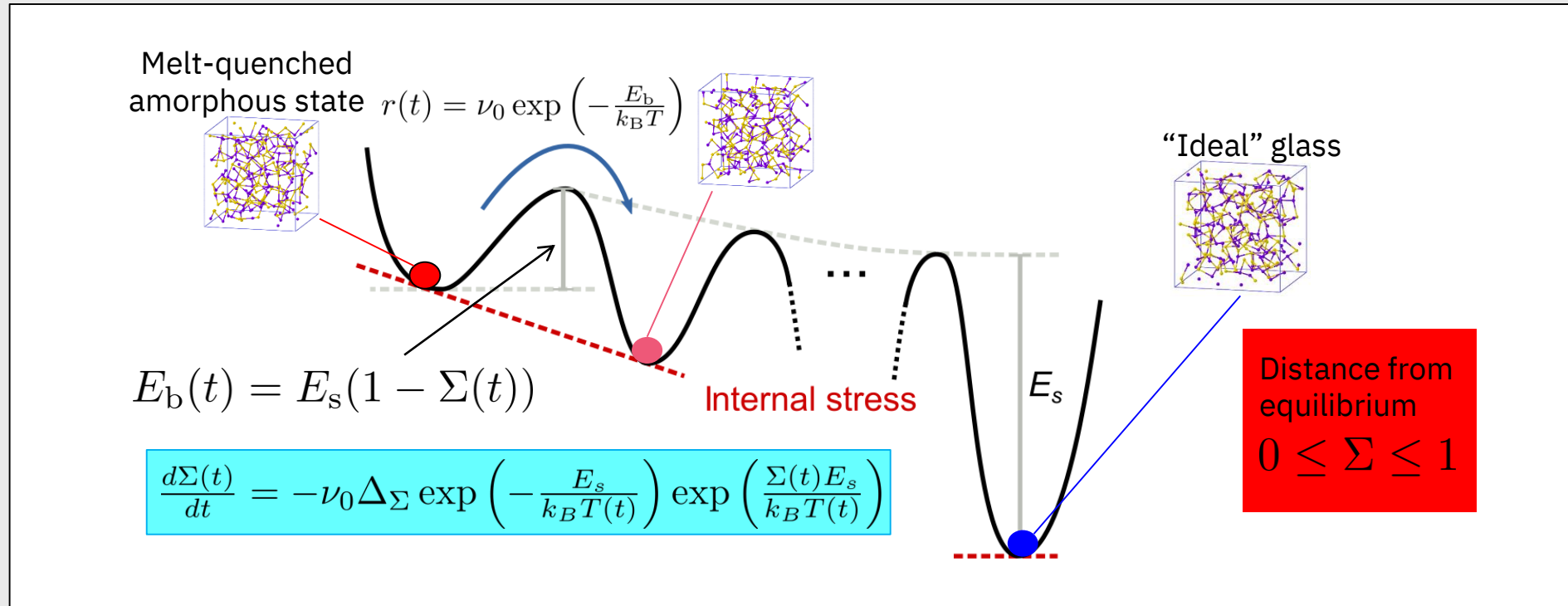
*Ielmini et al., IEEE Trans. Elec. Dev., 2009*  
*Rizzi et al., Appl. Phys. Lett., 2011*  
*Fantini et al., Appl. Phys. Lett., 2012*



Recent insights into the nature of the “ideal”  
glass state

*Raty et al., Nature Comm., 2015*  
*Gabardi et al., Phys. Rev. B, 2015*  
*Zipoli et al., Phys. Rev. B, 2016*

# The collective relaxation model



- The structure collectively rearranges whereby every local configuration is changed repeatedly to achieve an overall lower energy state
- The relaxation proceeds in a sequence of transitions between neighboring states
- Closer to the equilibrium the system is  $\rightarrow$  the higher the barrier for subsequent relaxation

**Sebastian et al., Proc. IRPS, 2015, Le Gallo et al., Proc. IRPS, 2016**

**Le Gallo et al., “Collective Structural Relaxation in Phase-Change Memory Devices”, Adv. Electr. Mat., 2018**

# The link between collective relaxation and electrical transport

## The relaxation dynamics @ constant temperature

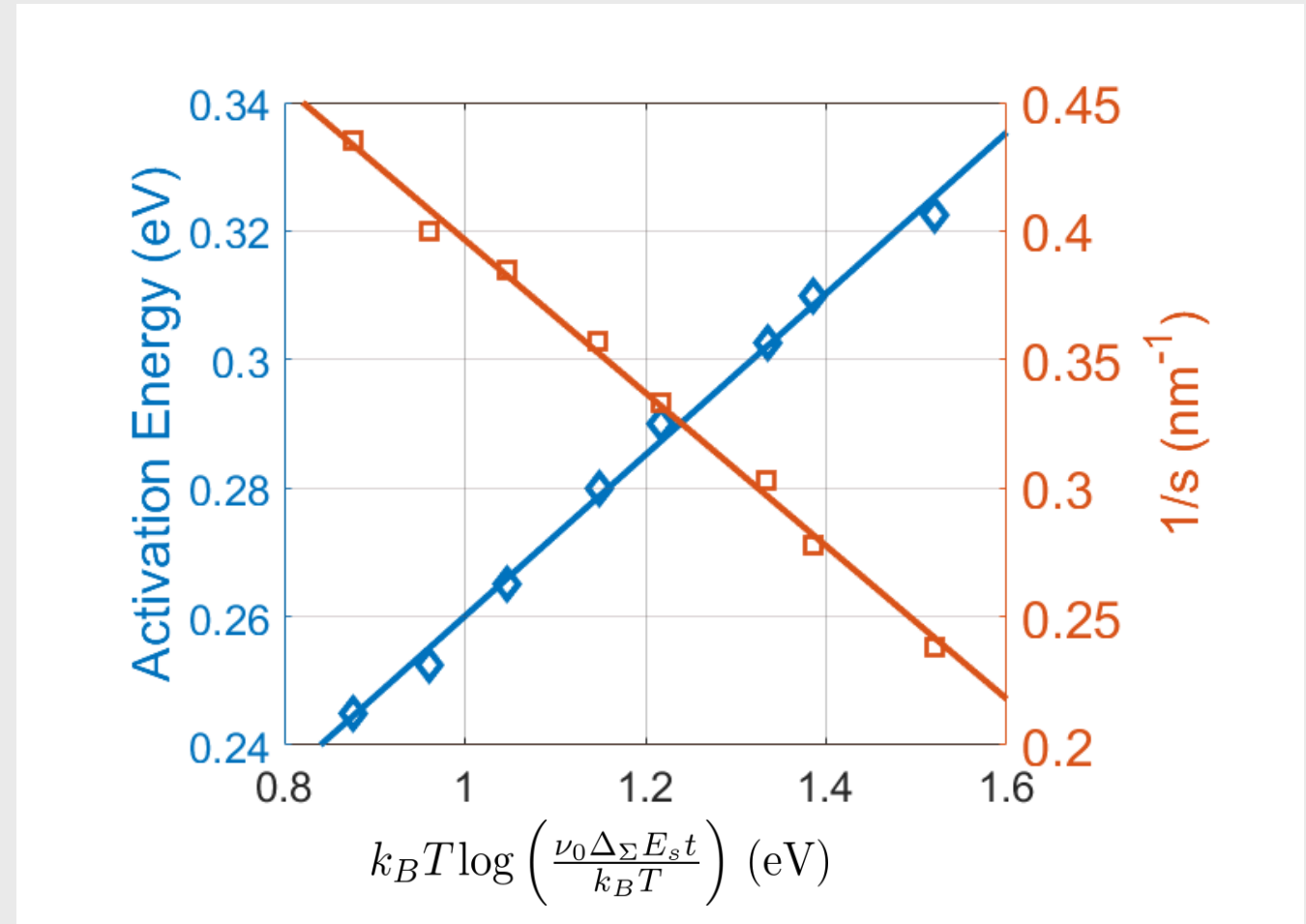
$$\Sigma(t) = -\frac{k_B T}{E_s} \log \left( \frac{t + \tau_0}{\tau_1} \right)$$
$$\tau_1 = (k_B T / \nu_0 \Delta \Sigma E_s) \exp \left( \frac{E_s}{k_B T} \right)$$
$$\tau_0 = \tau_1 \exp \left( -\frac{\Sigma_0 E_s}{k_B T} \right)$$

if  $t \gg \tau_0$

$$\Sigma(t) = 1 - \frac{k_B T}{E_s} \log \left( \frac{\nu_0 \Delta \Sigma E_s t}{k_B T} \right)$$

$$\Delta \Sigma \propto k_B T \log \left( \frac{\nu_0 \Delta \Sigma E_s t}{k_B T} \right)$$

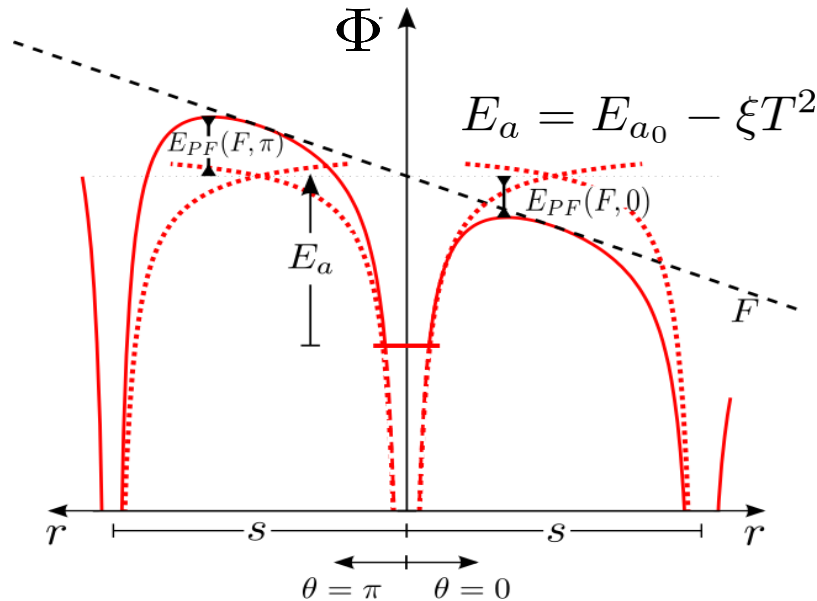
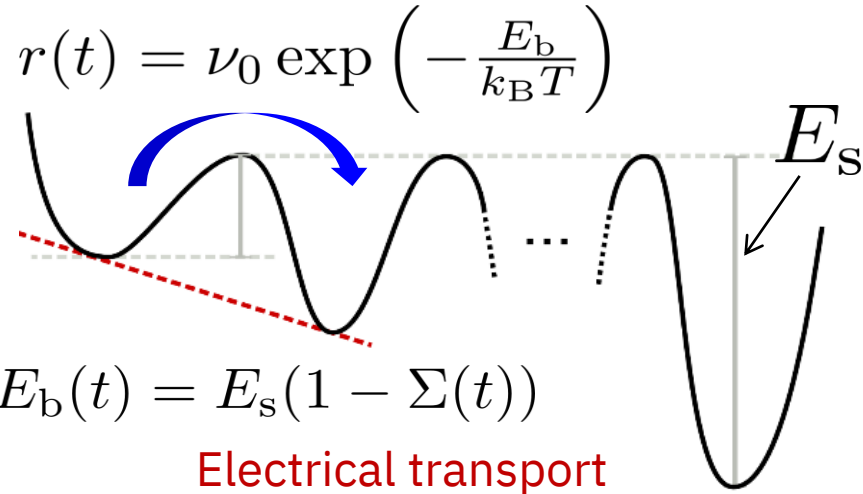
“Relaxation Energy”



Remarkable linear dependence of  $E_a$  and  $1/s$  on the relaxation energy

# Collective relaxation and electrical transport: The Link

## Structural relaxation



Amorphous state created with an initial distance from the “ideal glass”

$$\Sigma_0 = \Sigma(0)$$

Dynamics of structural relaxation

$$\frac{d\Sigma(t)}{dt} = -\nu_0 \Delta_\Sigma \exp\left(-\frac{E_s}{k_B T(t)}\right) \exp\left(\frac{\Sigma(t) E_s}{k_B T(t)}\right)$$

$$E_{a0}(t) = E_{a0}^{max} - \alpha \Sigma(t)$$

$$s(t) = s_0 / \Sigma(t)$$

Link to electrical transport

$$E_a(t) = E_{a0}(t) - \xi T^2$$

$$\Phi(r, \theta, F) = -eFr \cos(\theta) - \frac{\beta^2}{4e} \left( \frac{1}{r} + \frac{1}{s-r} \right) + \frac{\beta^2}{es}$$

$$E_{PF}(F, \theta) = -\max_r V(r, \theta, F)$$

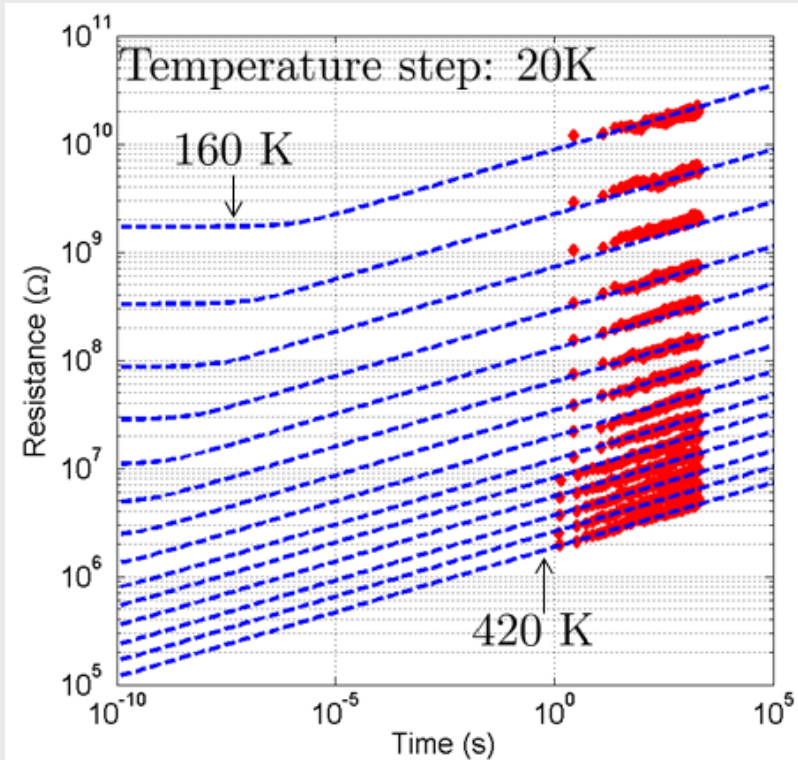
$$n(F) = \frac{K}{4\pi} \int_0^\pi \exp\left(-\frac{E_a - E_{PF}(F, \theta)}{k_B T}\right) 2\pi \sin(\theta) d\theta$$

$$\sigma(F) = e\mu n(F)$$

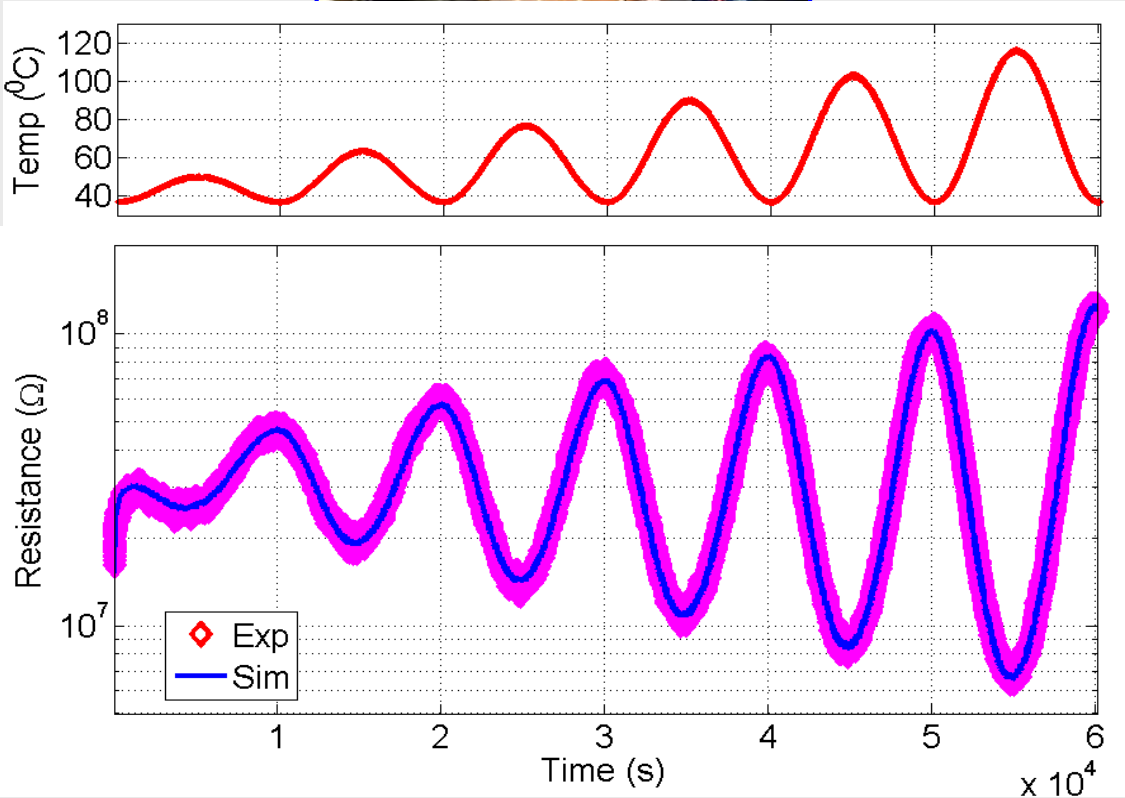
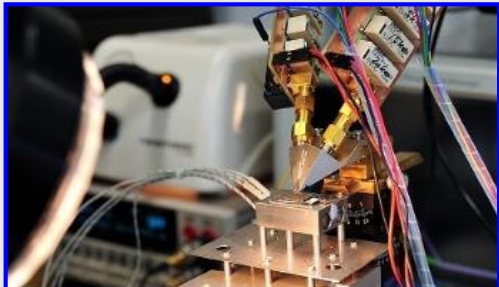


# Model validation

## Constant temperature over a wide range



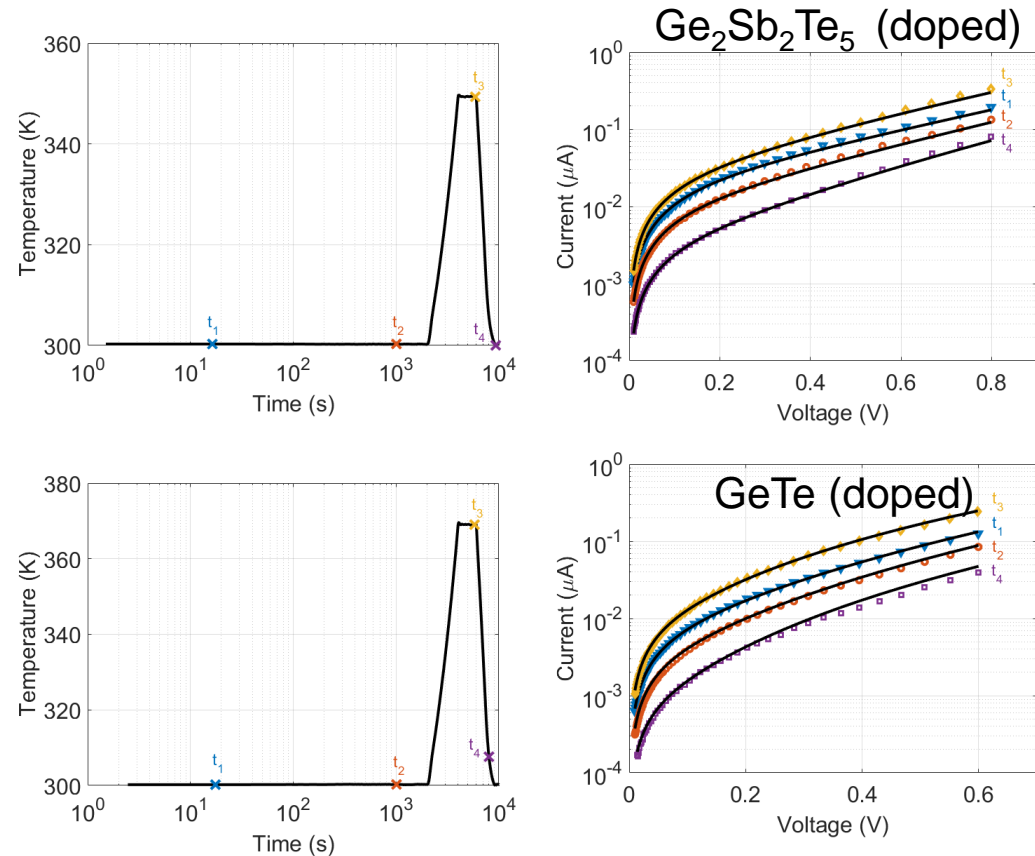
## Time-varying temperature



# Model validation: $\text{Ge}_2\text{Sb}_2\text{Te}_5$ and GeTe

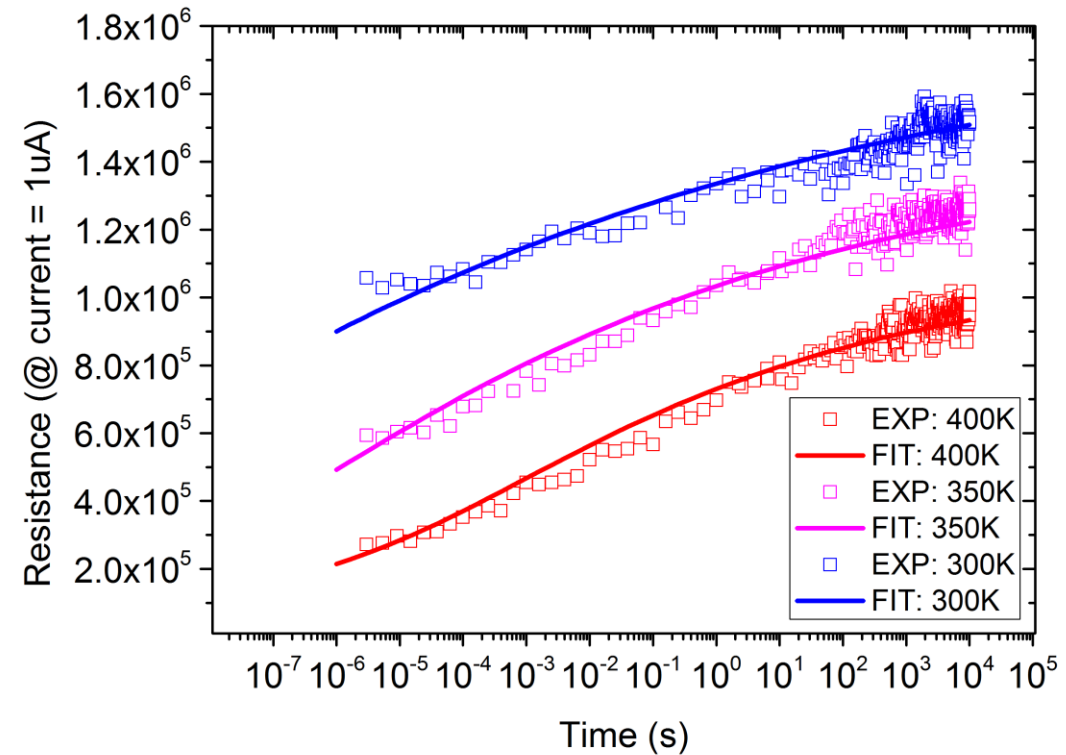
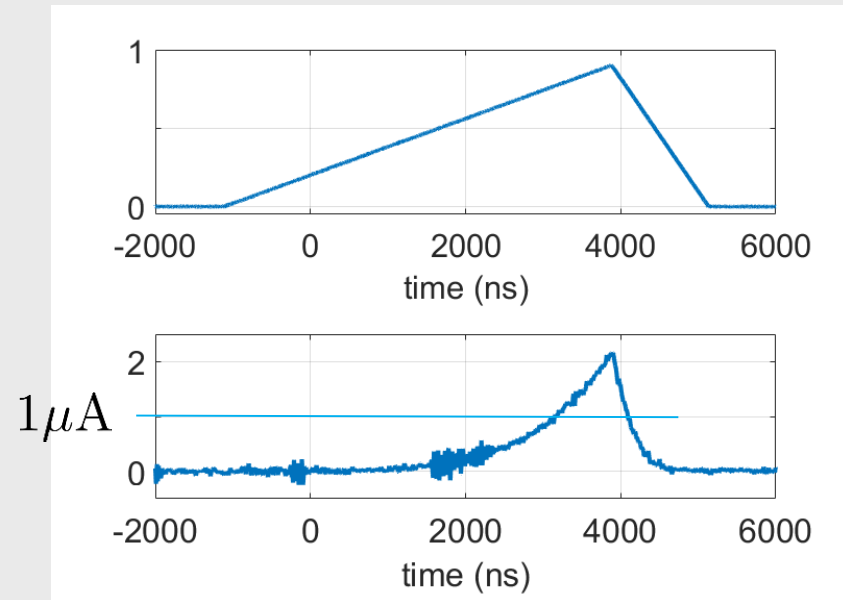
- Time-temperature profile applied to the PCM cell after RESET
- I-V curves accurately predicted by the drift model
- Clear increase of  $\log(I)$  vs  $V$  slope  $\rightarrow$  signature of increasing  $s$

Indicative of the universality  
of the drift model



# Model validation: High field resistance

**Fast measurement of high field resistance**

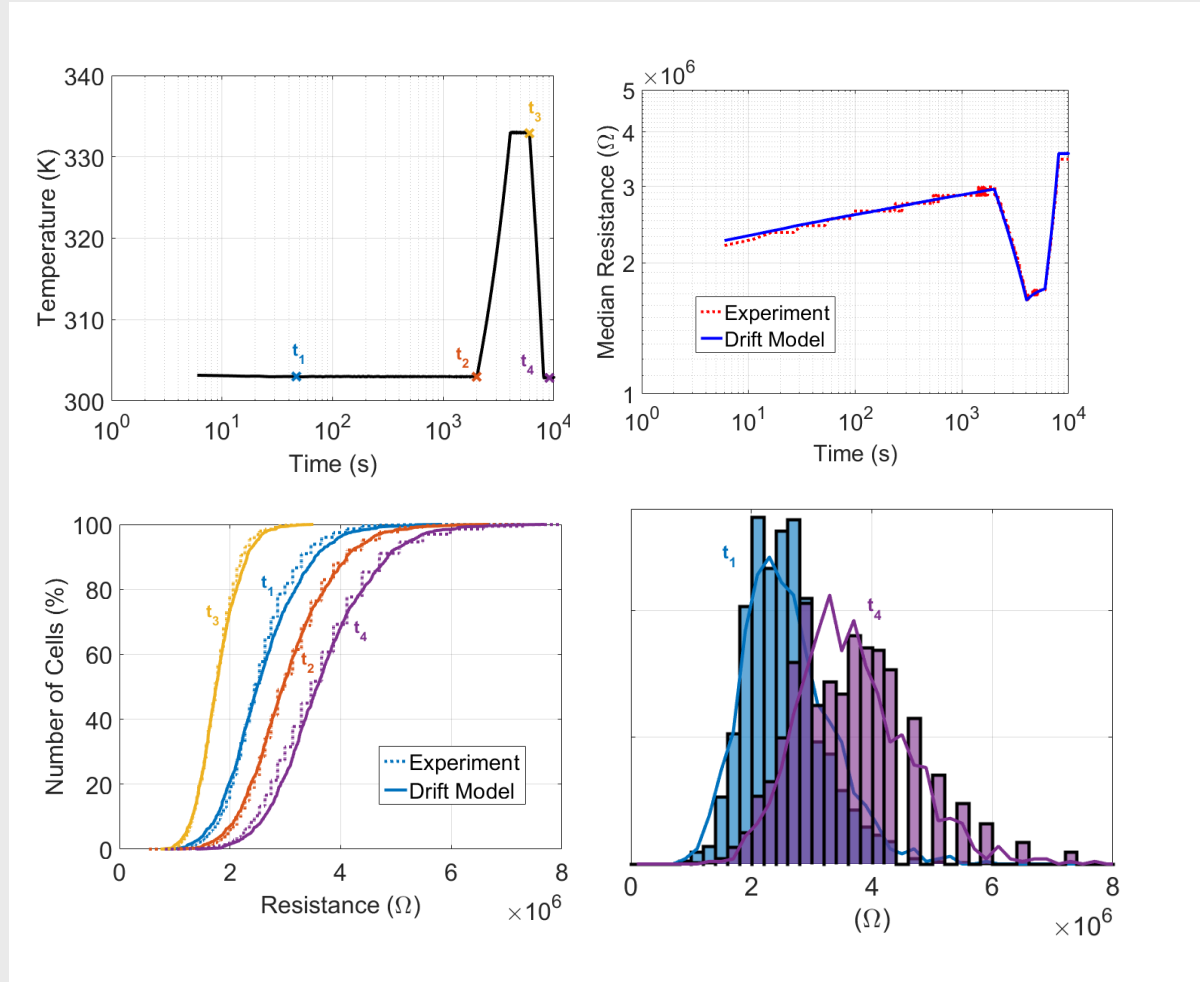


Model valid over 10 orders of magnitude in time

# Model validation: Array of 4k cells

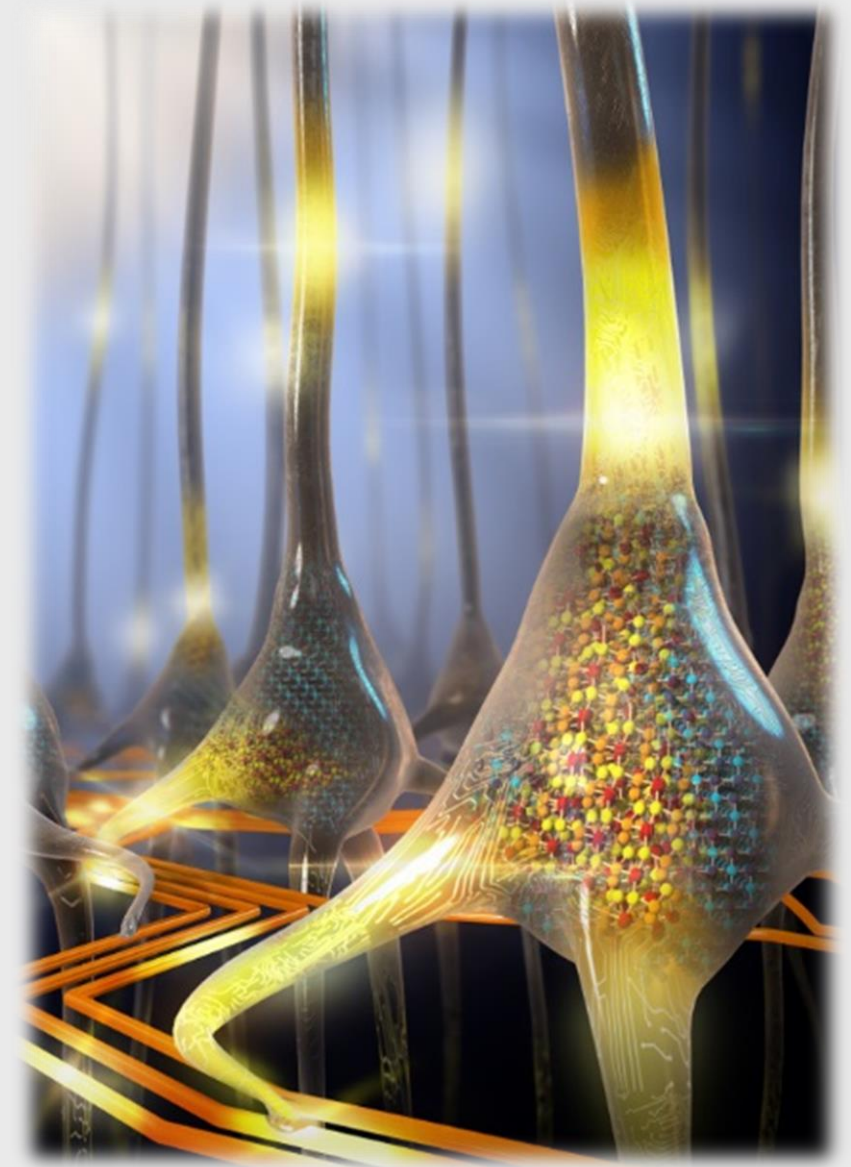
- A prototype PCM chip with addressing, readout and programming circuitry
- Measurements on a array of 4k cells undergoing changes in ambient temperature
- Model captures the evolution of the median resistance value
- Significant drift variability leads to a wide distribution of resistances even after short period of time
- Distributions broaden with increasing relaxation

Distributions and their broadening captured by the model



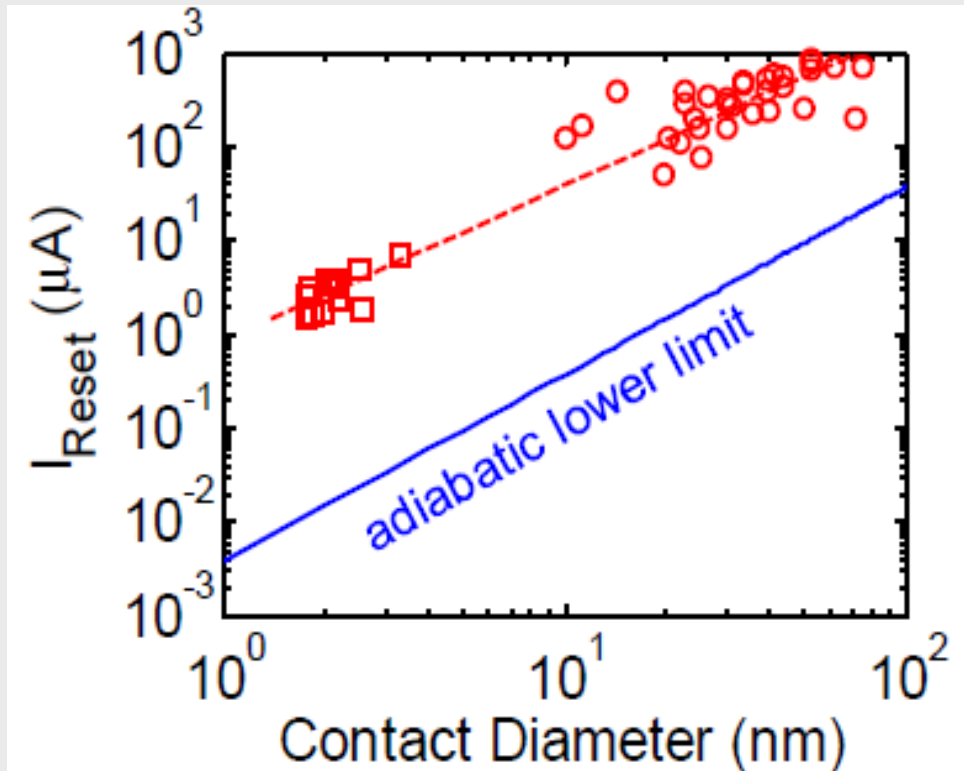
# Outline: Part I

- Introduction to PCM
- PCM device physics
  - ✓ **Electrical system**
    - Subthreshold electrical transport
    - Threshold switching
  - ✓ **Thermal system**
  - ✓ **Structural dynamics**
    - Melt-quench process
    - Crystallization
    - Structural relaxation
- **Key challenges and device-level advances**
  - ✓ Projected PCM
  - ✓ Single-elemental PCM

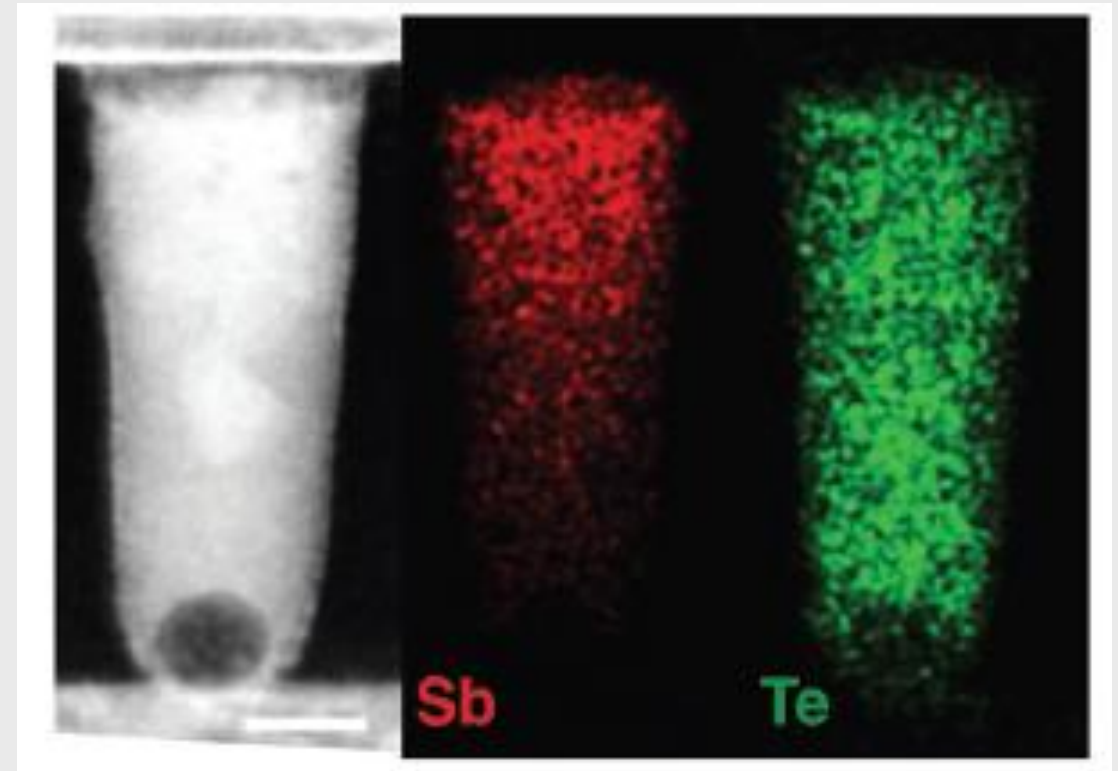




# Challenge 1: Programming energy



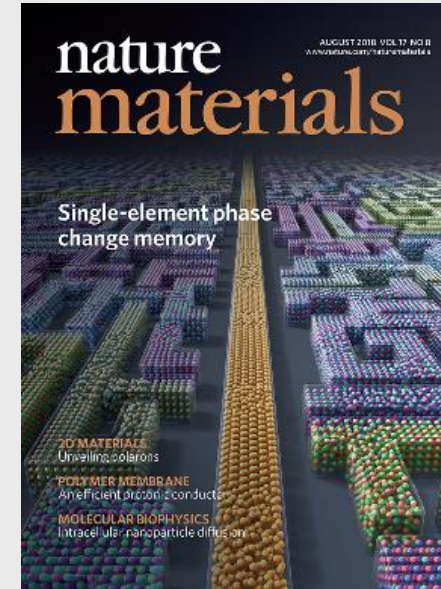
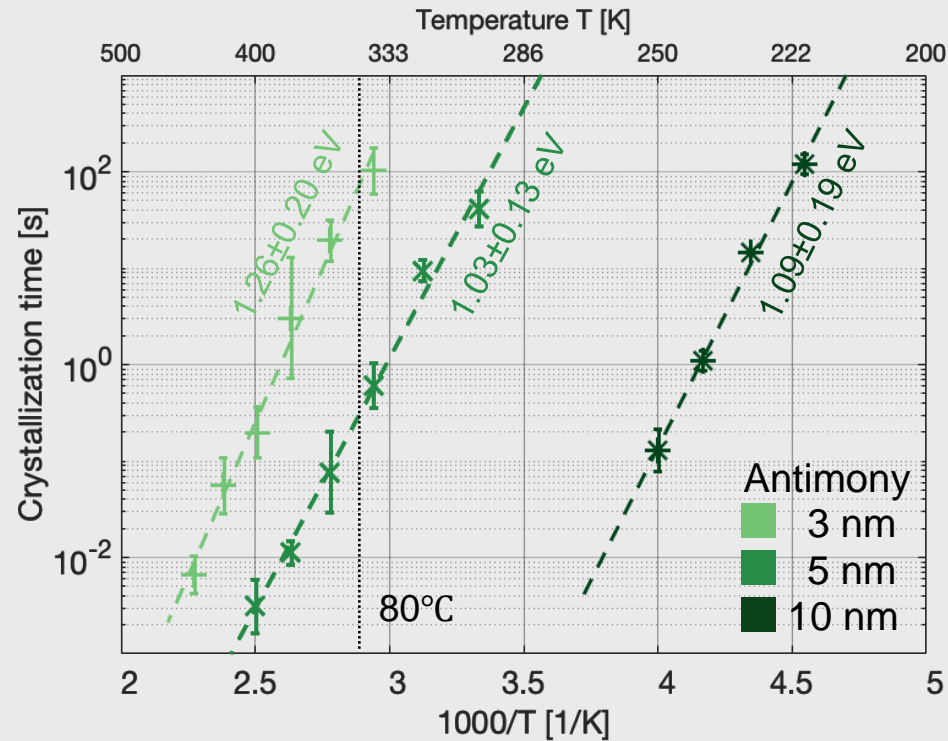
*Xiong et al., IEDM, 2016*



*Xie et al., Adv. Mat., 2018*

- Need to reduce the volume of material that needs to be melted
- Scaling to small dimensions severely hampered by fabrication challenges and elemental segregation

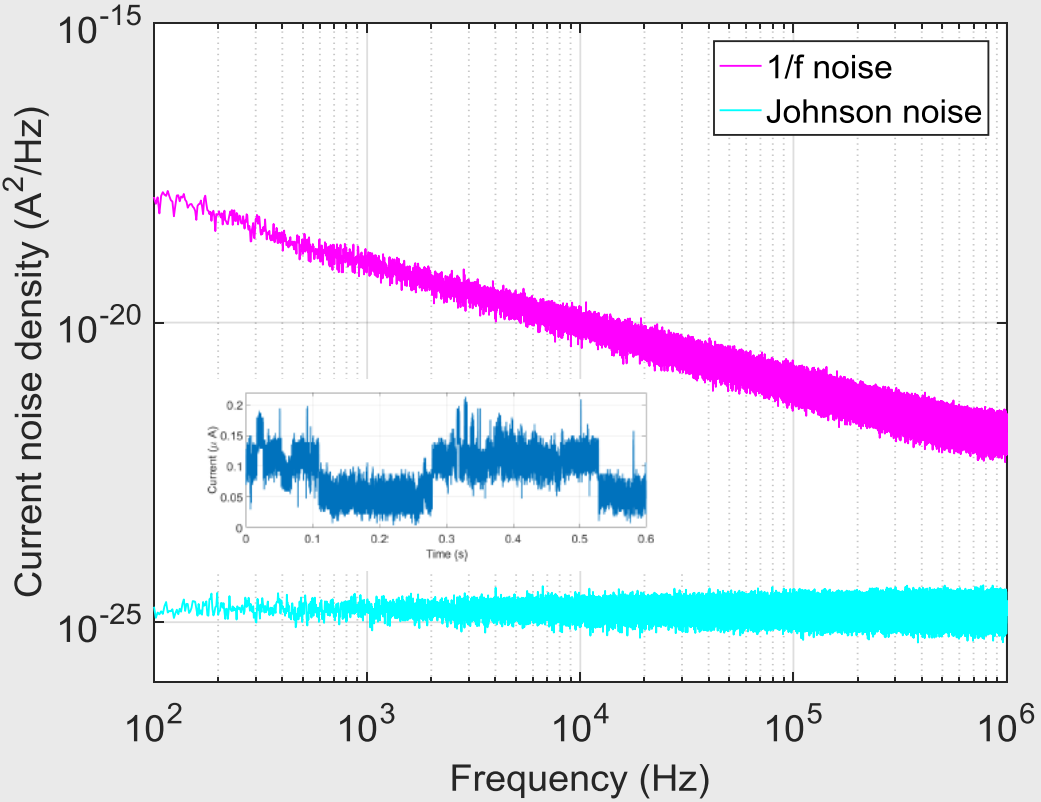
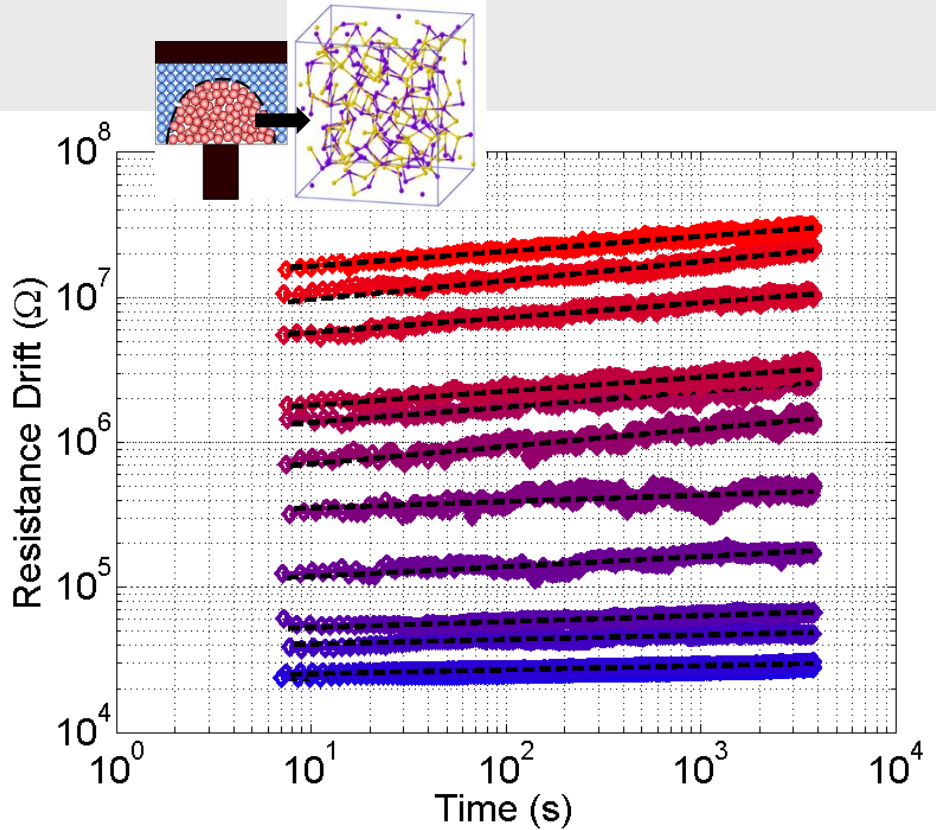
# Single element phase-change memory



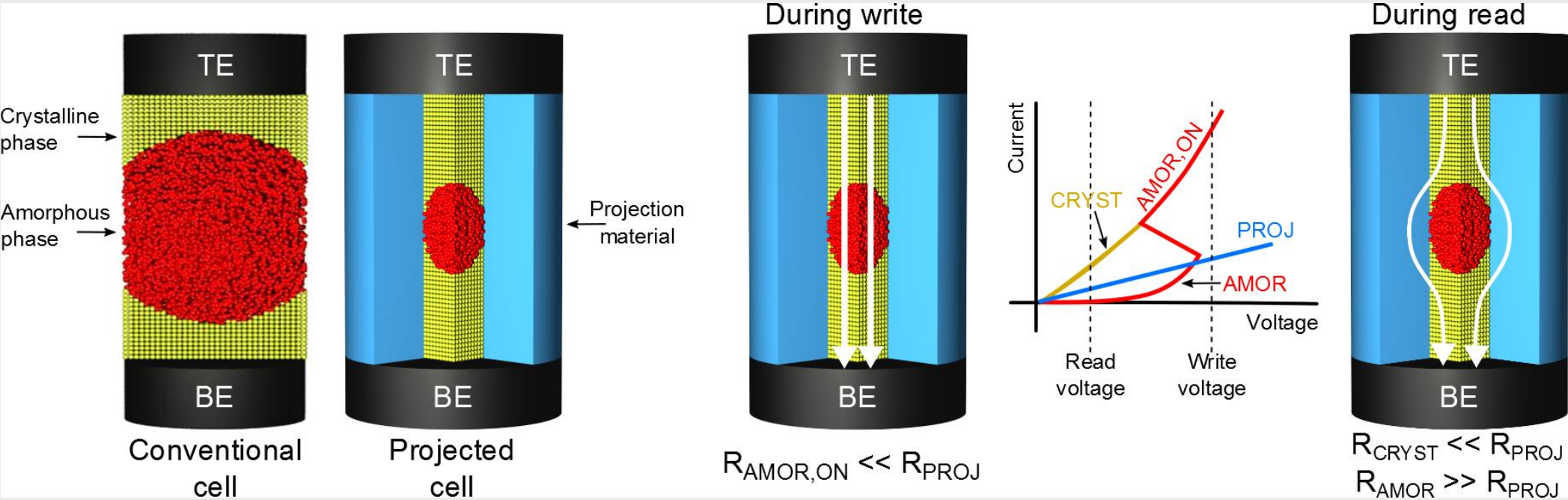
- **Single elemental Sb** can be used for PCM when confined to nanoscale dimensions!
- Directs research away from materials optimization to nanoscale confinement!

*Salinga et al., Nature Materials, 2018*

# Challenge 2: Drift and noise



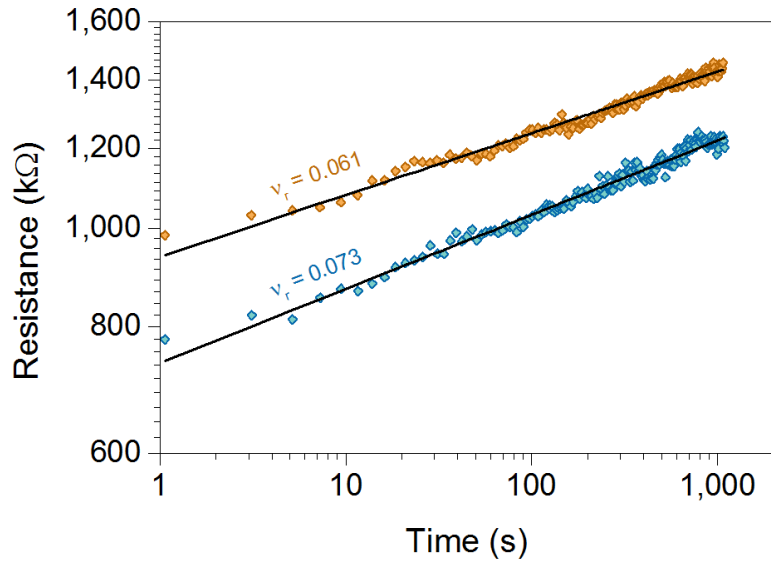
# Projected phase-change memory



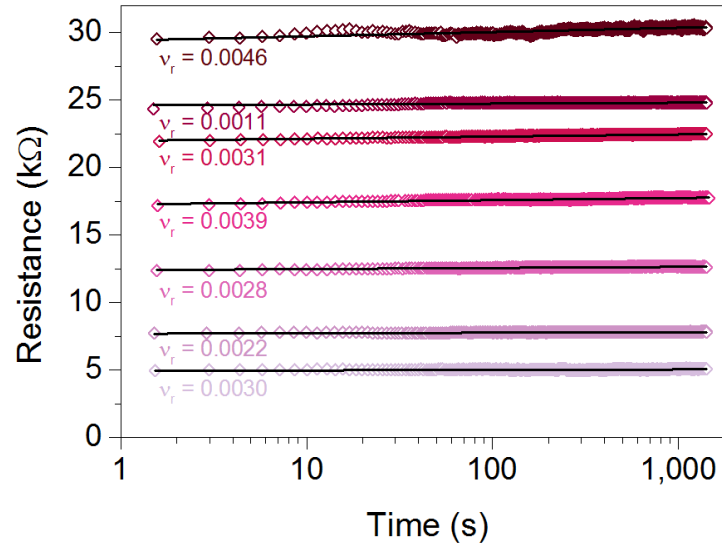
*Koelmans et al., Nature Comm., 2015*

# Projected phase-change memory

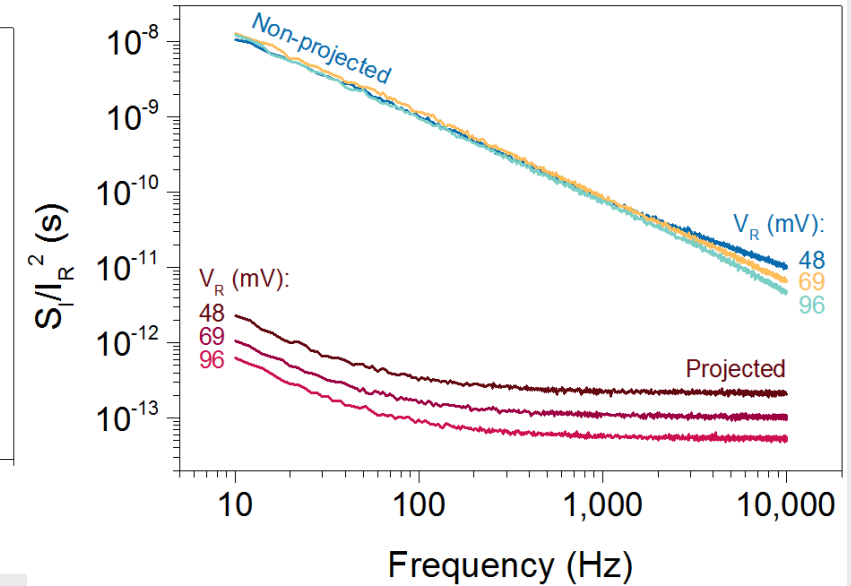
## Drift in un-projected devices



## Drift in projected devices



## Comparison of noise characteristics



- Near elimination of drift
- Substantial reduction in noise



# Summary

- **PCM is arguably the most advanced resistive memory technology**
  - ✓ The primary applications are foreseen in the domains of storage class memory and non-von Neumann computing
- **PCM device characteristics can be viewed as a intricate feedback interconnection of electrical, thermal and structural dynamics**
  - ✓ The electrical system
    - Subthreshold transport attributed to trap-limited band transport with PF emission
    - Threshold switching likely to arise from thermal feedback (at least in PCM devices)
  - ✓ Thermal system
    - Large temperature gradients within the devices → Thermoelectric effects need to be considered
  - ✓ Structural dynamics
    - Amorphization arises via melt-quench process
    - Crystal growth has a strong temperature dependence
    - Structural relaxation of the amorphous phase can be described via a collective relaxation model
- **Key challenges: The relatively large programming current and resistance fluctuations**
  - ✓ Single elemental PCM a promising avenue towards ultra-scaled PCM devices
  - ✓ Projected PCM, a device-level approach to countering resistance variations